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Grid-Voltage Sensorless Control of a Converter Under Unbalanced Conditions: On the Design of a State Observer

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Abstract—This paper deals with grid-voltage sensorless synchronization and control under unbalanced grid conditions. A three-phase grid-connected converter equipped with an LCL filter is considered, and no other signals than the converter currents and the DC-link voltage are measured for control. An augmented adaptive state observer is proposed for estimation of the positive- and negative-sequence components of the grid voltage. The proposed observer is tested as a part of a sensorless control system. Experimental results show that the proposed method works well even in highly unbalanced grid conditions.

Index Terms—Active front-end rectifier, grid-voltage sag, LCL filter, line-voltage sensorless, small-signal linearization.

I. INTRODUCTION

Grid-voltage sensorless control has been an interesting research topic during the last decades [1]–[19]. Replacing AC-voltage sensors with estimation is a relevant option, e.g., in pulse-width-modulated rectifiers [1]–[6], since reducing the number of the sensors can improve the reliability against electrical noise and reduce costs of the system. Furthermore, the grid-voltage sensorless operation has been proposed for distributed generation [8], [9]. Sensorless control, in parallel with the conventional operation, could increase reliability of a converter in renewable energy production, since a failure in measurement sensors, wires, or interfaces does not necessarily mean the end of production. In addition, a grid-voltage estimator might be useful in the islanding detection [11].

In the distributed generation, dynamic grid support, such as fault ride-through, is often required in high- and medium-voltage grids, but might be required in low-voltage grids as well in the future [20]. Furthermore, the converters may operate, at least temporarily, in unbalanced grid conditions [21], [22]. With careful design, unbalanced conditions can be also handled in grid-voltage sensorless control [4], [6], [9], [12]–[15], [19].

An LCL filter between the converter and the grid is an attractive option for filtering out the switching harmonics, because of its higher attenuation per volume or cost in comparison with the conventional L filter. Grid-voltage sensorless methods for converters equipped with the LCL filter have been proposed in [5], [8]–[10], [13]–[19]. The sensorless operation can be based on instantaneous power theory [10], virtual flux models, [5], [14]–[16], and model-based observers [8], [13], [17]–[19]. However, only a few of these methods take unbalanced grid conditions into account [9], [13]–[15], [19].

For the unbalanced grid conditions, a direct power control method has been developed in [14], but additional capacitor current measurements are needed for the active damping of the LCL-filter resonance. Additional capacitor current measurements have been used in [13] as well, where the grid-voltage estimation method is based on an adaptive algorithm [8]. In [9], an adaptive neural-network estimator has been proposed for disturbance and grid-voltage estimation. However, expertise in neural networks is needed in order to implement this type of estimator and to guarantee its stability. In [19], an adaptive model-based state observer has been used in grid-voltage sensorless control, but only the positive-sequence component of the grid voltage is estimated. In [15], grid-voltage sensorless operation is based on a virtual-flux model using a second-order generalized integrator as a fundamental building block. The method in [15] is an extension for the estimator proposed in [12].

This work deals with grid-voltage sensorless synchronization and control of a grid-connected converter under unbalanced grid conditions. A converter equipped with an LCL filter is considered, and the only measured quantities for control are the converter currents and DC-link voltage. Main contributions of this work are: 1) an augmented adaptive state observer is proposed for estimation of the positive- and negative-sequence components of the grid voltage; 2) tuning expressions for the observer are derived based on a small-signal linearized model and direct pole placement; 3) the proposed observer is experimentally tested as a part of a grid-voltage sensorless control system in unbalanced grid conditions.

II. SYSTEM MODEL

The equivalent circuit model of the LCL filter between the converter and the grid is shown in Fig. 1(a). The converter voltage is \( u^s_c \) and the converter current is \( i^s_c \), where the superscript \( s \) indicates stationary coordinates. The voltage across the filter capacitor is \( u^s_f \), the grid voltage is \( u^s_g \), and the grid current is \( i^s_g \). Complex-valued, matrix, and vector quantities are marked with boldface symbols.

In the case of unbalanced grid conditions (e.g. during a single- or two-phase fault), the grid-voltage vector

\[
\mathbf{u}^s_g(t) = \mathbf{u}^s_{g+}(t) + \mathbf{u}^s_{g-}(t) = e^{j\theta_{g+}(t)} u^s_{g+} + e^{j\theta_{g-}(t)} u^s_{g-} \quad (1)
\]
has the positive-sequence and negative-sequence components, \( u_{g+}^a \) and \( u_{g-}^a \), respectively. The positive-sequence magnitude is \( u_{g+} \) and the angle is

\[
\vartheta_{g+}(t) = \int \omega_{g+} \, dt
\]

where \( \omega_{g+} \) is the fundamental angular frequency of the grid voltage. For the negative-sequence component, \( u_{g-} \) is the magnitude, \( \vartheta_{g-} = -\vartheta_{g+} + \phi_{g-} \) is the angle, and \( \phi_{g-} \) is the phase shift.

A state-space model is formed from the equivalent circuit of the LCL filter [Fig. 1(a)]. The state vector is selected as \( x = [i_c^a, u_f^a, i_f^a]^T \). Moreover, the model is written in positive-sequence coordinates (marked with the superscript \( p \)). In these coordinates, the state vector is \( x^p = e^{-j\vartheta_{g+a}} x^a \), and other vector quantities are transformed in similar fashion. When the modeling principles presented in [23] are followed, the dynamics of the converter current \( i_c^a \) become

\[
x^p(k+1) = \Phi x^p(k) + \Gamma_c u_c^p(k) + \Gamma_g u_g^p(k) + \Gamma_{g+} u_{g+}^p(k) + \Gamma_{g-} u_{g-}^p(k)
\]

\[
i_c^p(k) = C_c x^p(k)
\]

where \( C_c = [1, 0, 0] \) is the output vector, \( \Phi \) is the state-transition matrix, \( \Gamma_c \) is the input vector for the converter voltage, and \( \Gamma_g \), \( \Gamma_g^+ \), and \( \Gamma_g^- \) are the input vectors for the positive- and negative-sequence components of the grid voltage, respectively. Symbolic expressions for calculating the model matrices are given in Appendix A.

In the discrete-time domain, the positive-sequence angle (2) is

\[
\vartheta_{g+}(k+1) = \vartheta_{g+}(k) + T_s \omega_{g+}
\]

where \( T_s \) is the sampling time. Furthermore, in positive-sequence coordinates, the rotating negative-sequence component satisfies the difference equation

\[
u_{g-}^p(k+1) = e^{-2j\omega_{g+} T_s} u_{g-}^p(k)
\]

For the purposes of disturbance estimation, the system model (3) is augmented with the negative sequence (5) as a disturbance state. The augmented state vector is \( x_0^p = [i_c^p, u_f^p, i_f^p, u_{g+}^p, u_{g-}^p]^T \), and the augmented system model is

\[
x_0^p(k+1) = \begin{bmatrix} \Phi & \Gamma_g^+ & \Gamma_g^- \\ 0 & e^{-2j\omega_{g+} T_s} & 0 \end{bmatrix} x_0^p(k) + \begin{bmatrix} \Gamma_g^+ \\ 0 \\ \Gamma_g^- \\ 0 \end{bmatrix} u_{g+}^p(k) + \begin{bmatrix} \Gamma_g^+ \\ 0 \\ \Gamma_g^- \\ 0 \end{bmatrix} u_{g-}^p(k)
\]

\[
i_c^p(k) = \begin{bmatrix} C_c \\ 0 \end{bmatrix} x_0^p(k)
\]

III. AUGMENTED ADAPTIVE OBSERVER

The proposed augmented adaptive observer is a part of the sensorless control system shown in Fig. 1(b). The structure of the proposed observer is shown in Fig. 2. While the augmented state vector \( x_0^p \) is estimated based on the model (6), the angular frequency \( \omega_{g+} \), positive-sequence angle \( \vartheta_{g+} \), and positive-sequence magnitude \( u_{g+} \) are estimated using adaptation mechanisms that are explained later.
The observer operates in estimated positive-sequence coordinates (marked without any superscripts). According to (4), the positive-sequence angle estimator is

$$\hat{\theta}_g^+ (k+1) = \hat{\theta}_g^+ (k) + T_s \hat{\omega}_g^+$$

(7)

The estimated positive-sequence coordinate system is generally different from the actual positive-sequence coordinate system due to the possible estimation error $\hat{\theta}_g^+ = \theta_g^+ - \hat{\theta}_g^+$, as illustrated in Fig. 3.

In estimated coordinates, the augmented state vector is $\tilde{x}_a = [x^T, u_{g-1}, u_g]^T$, and according to (6), the augmented state observer is

$$\hat{x}_a (k+1) = \hat{\Phi}_a \hat{x}_a (k) + \hat{\Gamma}_{ca} u_c (k) + \hat{\Gamma}_{ga} \hat{u}_g^+ (k) + K_o [\dot{i}_c (k) - \hat{i}_c (k)]$$

(8)

where $\hat{\Phi}_a$, $\hat{\Gamma}_{ca}$, and $\hat{\Gamma}_{ga}$ are the adaptive model matrices that are obtained by replacing the actual angular frequency $\omega_g^+$ with the estimated angular frequency $\hat{\omega}_g^+$ in (6). Furthermore, $K_o$ is the observer gain vector and $\hat{i}_c (k) = C_a \tilde{x}_a (k)$.

### A. Estimation-Error Dynamics

The estimation-error dynamics, derived in the following, play a key role in the tuning of the proposed observer. First, from (4) and (7), the dynamics of the angle estimation error become

$$\hat{\theta}_g^+ (k+1) = \hat{\theta}_g^+ (k) + T_s \hat{\omega}_g^+$$

(9)

where $\hat{\omega}_g^+ = \omega_g^+ - \hat{\omega}_g^+$ is the estimation error of the angular frequency. The system model (6) and the observer (8) are in different coordinates. For calculation of the state estimation error, the system model is transformed into estimated positive-sequence coordinates, where the observer operates (cf. Fig. 3).

The transformation of the state vector is

$$x^p_a (k) = e^{-j \hat{\theta}_a^+ (k)} x_a (k)$$

(10)

and $u_{g}^p$ and $\dot{i}_c^p$ are transformed similarly. The state vector at $k+1$ is transformed as $x^p_a (k+1) = e^{-j \hat{\theta}_a^+ (k+1)} x_a (k+1)$. Together with (9), the transformed system becomes

$$x_a (k+1) = e^{j T_s \hat{\omega}_g^+ (k)} \Phi_a x_a (k) + e^{j T_s \hat{\omega}_g^+ (k)} \Gamma_{ca} u_c (k) + e^{j \hat{\theta}_a^+ (k+1)} \Gamma_{ga} u_g^+ (k)$$

$$\dot{i}_c (k) = C_a x_a (k)$$

(11)

The estimation error of the augmented state vector is $\tilde{x}_a = x_a - \hat{x}_a$. From (8), (9), and (11), the estimation-error dynamics become

$$\hat{x}_a (k+1) = (\hat{\Phi}_a - K_o C_a) \hat{x}_a (k) + \left( e^{j T_s \hat{\omega}_g^+ (k)} \Phi_a - \hat{\Phi}_a \right) x_a (k) + \left( e^{j T_s \hat{\omega}_g^+ (k)} \Gamma_{ca} u_c (k) + \hat{\Gamma}_{ga} \hat{u}_g^+ (k) \right) + \left( e^{j \hat{\theta}_a^+ (k+1)} + e^{j \hat{\theta}_a^+ (k)} \right) \Gamma_{ga} \hat{u}_g^+ (k)$$

(12)

where $\hat{u}_g^+ = u_{g+} - \hat{u}_g^+$. The estimation-error dynamics (12) are linearized at an equilibrium point ($\tilde{x}_a = 0$, $u_{g+} = 0$, $\dot{\omega}_g^+ = 0$, $\dot{\theta}_g^+ = 0$), where the steady-state estimation errors are zero. Close to the equilibrium point, the dynamics are described by a linear state equation

$$\hat{x}_a (k+1) = (\Phi_a - K_o C_a) \hat{x}_a (k) + \Gamma_{ga} \hat{u}_g^+ (k)$$

$$+ j \Gamma_{ga} u_{g+} \hat{\theta}_g^+ (k) + \Gamma_{\omega} \hat{\omega}_g^+ (k)$$

(13)

where $\Gamma_{\omega}$ is the input vector for the angular-frequency estimation error. These linearized dynamics are derived in Appendix B.

In order to simplify the tuning of the adaptive observer, the linearized parameters can be further simplified by approximating $\Gamma_{\omega} \approx 0$. The approximation is reasonable, since the impact of the input $\hat{\omega}_g^+$ on $\tilde{x}_a$ in (8) is small in comparison with the other inputs ($\hat{u}_g^+$ and $\hat{\theta}_g^+$). The simplified dynamics are

$$\hat{x}_a (k+1) \approx (\Phi_a - K_o C_a) \hat{x}_a (k) + \Gamma_{ga} \hat{u}_g^+ (k) + j \Gamma_{ga} u_{g+} \hat{\theta}_g^+ (k)$$

(14)

The converter-current estimation error $\hat{i}_c = C_a \tilde{x}_a$ is the input for the positive-sequence frequency and magnitude estimators (cf. Fig. 2). In order to tune the estimators, the impact of $\hat{\theta}_g^+$ and $\hat{\omega}_g^+$ on $\hat{i}_c$ is studied. From (14), the pulse-transfer function from $\hat{u}_g^+$ to $\hat{i}_c$ becomes

$$G_{iu} (z) = \frac{\hat{i}_c (z)}{\hat{u}_g^+ (z)} = C_a (z I - \Phi_a + K_o C_a)^{-1} \Gamma_{ga} u_{g+}$$

(15)

1The impact of the inputs at $k$ on $\tilde{x}_a (k+1)$ can be evaluated by calculating the gains $\| \Gamma_{\omega} \|$ and $\| \Gamma_{\omega} \|$. Starting from the equilibrium $\tilde{x}_a (k) = 0$, the change $\| \tilde{x}_a (k+1) \|$ in the states caused by the change equal to unity in $\hat{\omega}_g^+$ is $\| \Gamma_{\omega} \|$. For example, with the parameters of Table 1 and symmetrical grid-voltage conditions, the gain $\| \Gamma_{\omega} \| = 0.01$ p.u., whereas the gain $\| \Gamma_{\omega} \| = \| \Gamma_{\omega} u_{g+} \| = 0.56$ p.u.
where
\[ a(z) = \text{det}(zI - \Phi_a + K_a C_a) \]  
and \( b(z) \) is the numerator of the transfer function. Similarly, the pulse-transfer function from \( \hat{\tilde{\epsilon}}_g \) to \( \hat{\tilde{\epsilon}}_c \) becomes
\[ G_{g\hat{\tilde{\epsilon}}}(z) = \frac{\tilde{\epsilon}_c(z)}{\hat{\tilde{\epsilon}}_g(z)} = \frac{ju_{g+0} \cdot b(z)}{a(z)} \]  

C. Adaptation Laws

The proposed algorithms for estimating the positive-sequence angle \( \hat{\omega}_{g+} \), magnitude \( \hat{u}_{g+} \), and frequency \( \hat{\omega}_{g+} \) are based on the small-signal linearized model (14). The estimator for the positive-sequence angle is given in (7), and the estimator for the positive-sequence magnitude is
\[ \hat{u}_{g+}(z) = \frac{k_{i,\text{u}}}{z - 1} \text{Re} \left\{ \frac{a_1}{b_1} e^{j\phi \gamma} \tilde{\epsilon}_c(z) \right\} \]  

where \( k_{i,\text{u}} \) is the gain, and \( \phi, a_1, \) and \( b_1 \) are explained at the end of this subsection. The estimator for the angular frequency is
\[ \hat{\omega}_{g+}(z) = \frac{1}{u_{g+0}} \left( k_{p,\omega} + k_{i,\omega} \right) \text{Im} \left\{ \frac{a_1}{b_1} e^{j\phi \gamma} \tilde{\epsilon}_c(z) \right\} \]  

where \( u_{g+0} \) is the positive-sequence voltage at the operating point, and \( k_{p,\omega} \) and \( k_{i,\omega} \) are the estimator gains. Furthermore, the integral part of (19) provides a naturally filtered frequency estimate, e.g., for monitoring purposes [19]
\[ \hat{\omega}_{\text{ref}}(z) = \frac{1}{u_{g+0}} k_{i,\omega} \text{Im} \left\{ \frac{a_1}{b_1} e^{j\phi \gamma} \tilde{\epsilon}_c(z) \right\} \]  

The constants \( \phi, a_1, \) and \( b_1 \) are obtained from the quasi-steady-state analysis of the linearized model, as explained in [19]. These constants describe the steady-state gain from \( u_{g+} \) and \( \hat{\omega}_{g+} \) to \( \tilde{\epsilon}_c \) as follows
\[ \hat{\tilde{\epsilon}}_c = e^{-j\theta} \left( \frac{b_1}{a_1} \right) \hat{u}_{g+} + j u_{g0} e^{-j\theta} \left( \frac{b_1}{a_1} \right) \hat{\omega}_{g+} \]  

and they are obtained by making \( z = 1 \) in (15) and (17). The constants are
\[ \phi = (3/2) \omega_{g+} T_s \]
\[ a_1 = \omega_{g+} C_l L_f \left( \omega_{g+}^2 - \omega_p^2 \right) \times (1 - \alpha_{o1})(1 - \alpha_{o2})(1 - \alpha_{o3})(1 - \alpha_{o4}) \]
\[ b_1 = (4 - 4e^{-2j\omega_{g+} T_s}) \sin(\omega_{g+} T_s/2) \times [\cos(\omega_{g+} T_s) - \cos(\omega_p T_s)] \]

where \( \omega_p = \sqrt{(L_f + L_{fg})/(C_l L_f L_{fg})} \) is the resonance frequency of the LCL filter and \( \alpha_{o1}, \alpha_{o2}, \alpha_{o3}, \) and \( \alpha_{o4} \) are the discrete-time poles (tuning parameters) of the state observer.

D. Tuning of the Observer

The proposed observer is tuned based on direct pole placement. Two sets of poles are placed: poles of the state observer and poles of the adaptation loops.

1) Poles of the State Observer: The gain \( K_o \) is calculated by selecting the pole locations of the augmented state observer. These poles are the roots of (16). Four complex poles \( \alpha_{o1} \ldots \alpha_{o4} \) can be placed resulting in the desired characteristic polynomial
\[ a(z) = (z - \alpha_{o1})(z - \alpha_{o2})(z - \alpha_{o3})(z - \alpha_{o4}) \]  

Analytical expressions for the observer gain \( K_o \) are calculated by equalizing (16) and (23). The resulting expressions are functions of \( \alpha_{o1} \ldots \alpha_{o4} \) and the elements of \( \Phi_a \).

In order to simplify the selection of the pole locations, the discrete-time poles are mapped via the continuous-time polynomials \( s^2 + 2\zeta_{\text{od}} \omega_{\text{od}} s + \omega_{\text{od}}^2 \) and \( s^2 + 2\zeta_{\text{or}} \omega_{\text{or}} s + \omega_{\text{or}}^2 \) resulting in
\[ \alpha_{o1,2} = \exp[-\zeta_{\text{od}} \pm j \sqrt{1 - \zeta_{\text{od}}^2} \omega_{\text{od}} T_s] \]
\[ \alpha_{o3,4} = \exp[-\zeta_{\text{or}} \pm j \sqrt{1 - \zeta_{\text{or}}^2} \omega_{\text{or}} T_s] \]  

where the natural frequencies (\( \omega_{\text{od}}, \omega_{\text{or}} \)) and damping ratios (\( \zeta_{\text{od}}, \zeta_{\text{or}} \)) determine the location of the poles. Following the direct pole-placement principle proposed in [23], the natural frequency \( \omega_{\text{od}} \) is here set twice as fast as the bandwidth of current control. The natural frequency \( \omega_{\text{or}} \) is set to the resonance frequency of the LCL filter, i.e., \( \omega_{\text{or}} = \omega_p \). It is to be noted that this selection is a practical example, and a designer has the full freedom to change the pole locations depending on the system specifications and conditions. For example, if the measured current \( \tilde{\epsilon}_c \) is noisy, the natural frequencies can be slowed down in order to increase measurement-noise rejection.

### Table I

<table>
<thead>
<tr>
<th>Param.</th>
<th>Value</th>
<th>Param.</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>( u_{g+} )</td>
<td>( \sqrt{2/3} \cdot 400 \text{ V (1 p.u.)} )</td>
<td>( L_{fc} )</td>
<td>3.3 mH (0.081 p.u.)</td>
</tr>
<tr>
<td>( \omega_{g+} )</td>
<td>( 2\pi \cdot 50 \text{ rad/s (1 p.u.)} )</td>
<td>( L_{fg} )</td>
<td>3.0 mH (0.074 p.u.)</td>
</tr>
<tr>
<td>( i_{\text{B}} )</td>
<td>( \sqrt{2} \cdot 18 \text{ A (1 p.u.)} )</td>
<td>( C_l )</td>
<td>8.8 ( \mu \text{F} ) (0.036 p.u.)</td>
</tr>
<tr>
<td>( f_{sw} )</td>
<td>4 kHz</td>
<td>( T_2 )</td>
<td>1/(2f_{sw}) = 120 ( \mu \text{s} )</td>
</tr>
</tbody>
</table>
2) Poles of the Adaptation Loops: Since the adaptation mechanisms are slow in comparison with the state observer, the quasi-steady-state approximation (21) of the linearized dynamics can be used for tuning of the positive-sequence magnitude, frequency, and angle estimators. The estimators (7), (18), (19) form loops together with (21), as shown in Fig. 4. The gains for the estimators are calculated by selecting the pole locations of these closed adaptation loops. The resulting gain of the positive-sequence magnitude estimator is

\[ k_{i,u} = 1 - \exp(-\omega_u T_s) \]  

where \( \omega_u \) is the natural frequency of the pole of the magnitude-estimation loop. Furthermore, the gains of the angle-estimation loops become

\[ k_{i,\omega} = \frac{2[1 - \exp(-\zeta_\omega \omega_w T_s)] \cos(\sqrt{1 - \zeta_\omega^2} \omega_w T_s)}{\zeta_\omega \omega_w T_s}, \quad k_{i,\phi} = [\exp(-2\zeta_\phi \omega_w T_s) - 1] T_s + k_{p,\omega} \]  

where \( \omega_w \) is the natural frequency and \( \zeta_\omega \) is the damping ratio of the complex pole pair of the angle-estimation loop.

The natural frequencies \( \omega_u \) and \( \omega_w \) and the damping ratio \( \zeta_\omega \) are the tuning parameters of the loops, and the natural frequencies can be thought as approximate bandwidths. It is worth noticing that the linearized system [(15) and (17)] has a zero (\( z = e^{-2j\omega_u T_s} \)) at the unit circle\(^2\) that limits the maximum values of \( \omega_u \) and \( \omega_w \).

E. Small-Signal Stability

The small-signal stability of the proposed observer is analyzed in the nominal operation point. In the stability analysis, \( \tilde{u}_{g+}(k) = u_{g+}(k) - \hat{u}_g(k), \tilde{\omega}_g(k) = \omega_g(k) - \hat{\omega}_g(k), \) and \( \tilde{\omega}_{g+}(k) = \omega_{g+}(k) - \hat{\omega}_{g+}(k), \) where \( u_{g+} \) and \( \omega_{g+} \) as external disturbances can be set to zero. Considering (9), (13), and (18)–(20), the small-signal model for the closed-loop system is obtained as follows

\[
\begin{align*}
\dot{\tilde{x}}_a(k + 1) &= (\Phi_a - K_a C_a)\tilde{x}_a(k) + \Gamma_g \tilde{u}_{g+}(k) + j\Gamma_g u_{g+} \tilde{\omega}_{g+}(k) \\
\hat{\omega}_{g+}(k + 1) &= \hat{\omega}_{g+}(k) - k_{i,\omega} \frac{a_1}{b_1} e^{j\phi} C_a \tilde{x}_a(k) \\
\tilde{\omega}_{g+}(k + 1) &= \tilde{\omega}_{g+}(k) + \frac{k_{i,\omega}}{u_{g+}} \frac{a_1}{b_1} e^{j\phi} \tilde{x}_a(k) \\
\hat{\vartheta}_{g+}(k + 1) &= \hat{\vartheta}_{g+}(k) + T_s \tilde{\omega}_{g+}(k)
\end{align*}
\]

where

\[
\tilde{\omega}_{g+}(k) = -\frac{k_{p,\omega}}{u_{g+}} \frac{a_1}{b_1} e^{j\phi} \hat{\vartheta}(k) + \tilde{\omega}_{g+}(k)
\]

A numerical example of the observer tuning is presented. The system parameters are given in Table I. The tuning parameters of the state observer are: \( \omega_{od} = 2\pi \cdot 1000 \) rad/s, \( \omega_{or} = \omega_p, \zeta_{od} = 0.9, \) and \( \zeta_{or} = 0.7. \) Fig. 5 shows the poles of the closed-loop system (27), when \( \omega_u \) is increased from 2\( \pi \cdot 5 \) rad/s to 2\( \pi \cdot 100 \) rad/s, \( \omega_w = \omega_u, \) and \( \zeta_w = 1. \) The adaptation mechanisms change the original location (marked with green *) of the state-observer poles determined by \( K_a. \) Two state observer poles (blue) drift outside of the stable region (unit circle) when \( \omega_u = \omega_w > 2\pi \cdot 65 \) rad/s. On the other hand, when \( \omega_u = \omega_w < 2\pi \cdot 35 \) rad/s the damping ratios of all the poles are greater than 0.4. This means, that the estimation-error dynamics has a reasonable damping. The slow-frequency poles of the adaptation loops, determined by \( \omega_u \) and \( \omega_w, \) are in red.

IV. EXPERIMENTAL RESULTS

The proposed adaptive observer was experimentally tested as a part of the grid-voltage sensorless control scheme shown in Fig. 1(b). The switching frequency of the 12.5-kVA, 400-V converter under test is 4 kHz. The current controller is a state-space controller [23], which is tuned to give an approximate closed-loop bandwidth of 500 Hz. The controller is augmented with a reduced-order generalized integrator at \(-2\omega_{g+}\) for regulating the negative-sequence current component. Estimated and filtered positive- and negative-sequence grid-voltage components are used in the current reference calculation. A band-stop filter at \( 2\omega_{g+}\) removes the second-harmonic ripple from the measured DC-link voltage in the voltage control. The system parameters are given in Table I.

Direct pole placement was used in the observer tuning. The tuning parameters of the state observer are: \( \omega_{od} = 2\pi \cdot 1000 \) rad/s, \( \omega_{or} = \omega_p, \zeta_{od} = 0.9, \) and \( \zeta_{or} = 0.7. \) The tuning parameters for the adaptation loops are: \( \omega_u = \omega_w = 2\pi \cdot 25 \) rad/s, and \( \zeta_w = 1. \) The tuning corresponds to the case in the stability analysis of Section III-E.

A. Unbalanced Grid-Voltage Dips

Figs. 6 and 7 show the measured waveforms when the following grid-voltage sequence was generated using a 50-kVA four-quadrant power supply (Regatron TopCon TC.ACS): 1)
Fig. 6. Measured waveforms in different grid-voltage conditions: (first) grid voltages; (second) grid currents; (third) estimated positive-sequence magnitude (solid) and actual magnitude of the grid-voltage vector (dashed); (fourth) estimated positive-sequence angle (solid) and actual angle of the grid-voltage vector (dashed); and (fifth) estimated negative-sequence components.

Fig. 7. Measured locus of the grid-voltage space vector $u_g = u_g\alpha + j u_g\beta$: (a) symmetrical conditions; (b) the measured vector rotates an elliptic orbit when the single-phase dip is applied; and (c) the rotation stops (the vector is pulsating) when the two-phase dip is applied. The non-circular rotation of the vector is clearly visible in the measured magnitude $u_g = |u_g^e|$ and angle $\vartheta_g = \angle u_g^e$ in Fig. 6.

B. Phase-Angle Jump and Frequency Steps

Fig. 8. Measured waveforms during the phase-angle jump of $-60^\circ$ and frequency changes: (first) space-vector components of the grid-voltage; (second) space-vector components of the grid current; (third) estimated positive-sequence angle (solid) and actual angle of the grid-voltage vector (dashed); (fourth) estimated angular frequencies $\hat{\omega}_{g+}$ and $\hat{\omega}_{g-}$, and the actual angular frequency of the grid-voltage vector; and (fifth) estimated positive-sequence magnitude (dashed) and negative-sequence components (solid).

symmetrical grid conditions; 2) single-phase voltage dip down to zero; 3) two-phase voltage dip down to zero; 4) symmetrical grid conditions. During the test sequence, the converter was loaded by another back-to-back connected converter such that the converter under test was supplying the power of 0.3 p.u. Fig. 6 shows measured magnitude $u_g = |u_g^e|$ and angle $\vartheta_g = \angle u_g^e$ (blue dashed lines) of the grid-voltage vector (1). The locus of the vector is illustrated in Fig. 7. Theoretically, the positive-sequence magnitude in the voltage vector is 2/3 p.u. and 1/3 p.u. during the single-phase and two-phase dips, respectively. The negative-sequence magnitudes are 1/3 p.u. during the both dips. As Fig. 6 shows, the positive-sequence magnitude estimate converges quickly and the estimate agrees with the theoretical values. The positive-sequence angle is correctly detected as well, and it is only slightly perturbed in transients. The negative-sequence estimate also converges quickly and agrees with the theory.

As Fig. 6 shows, the grid currents are balanced even under highly unbalanced grid voltages. At the beginning, the magnitudes of the grid currents are 0.3 p.u. because the power of 0.3 p.u. is transferred. During the grid-voltage dips, the magnitudes of the currents increase first to 0.5 p.u. and then to 1 p.u. in order to transfer the same power and to maintain the power balance in steady state. The voltage transients shown in the figure are challenging but the sensorless control system is operating very well and it is able to achieve steady state approximately within a grid-voltage cycle.
naturally filtered frequency estimate \( \hat{\omega}_{gf} \) (green dashed line) given by (20) is free of impulses. During the frequency steps, the estimate \( \hat{\omega}_{g+} \) converges rapidly whereas the estimate \( \hat{\omega}_{gf} \) converges slower but it is noise-free. Some cross coupling is present between the estimated quantities, but the steady state is achieved approximately in 30 ms. In any case, the proposed observer and control system adapt to the frequency, which is shown by the grid-current space-vector components \( i_{ga} \) and \( i_{gb} \). The frequency of the current components follows that of the voltage components \( u_{ga} \) and \( u_{gb} \).

C. Feeding Negative-Sequence Current

The estimated negative-sequence component \( \hat{u}_{g-} \) can be used in the current reference calculation, for example, in order to eliminate \( 2\omega_{g+} \) ripple from the active power \( p_g \) or reactive power \( q_g \). Fig. 9 shows measured voltages, currents, and powers when a negative-sequence current component is added in the reference current at \( t = 0.04 \) s. The converter was supplying the power of 0.4 p.u., when the grid voltages were unbalanced (\( u_{g+} = 0.7 \) p.u. and \( u_{g-} = 0.3 \) p.u.). The reference for the negative-sequence current component was calculated using the estimated and filtered negative-sequence voltage component, and the reference was selected to eliminate the active-power ripple. As the figure shows, the ripple is reduced in the active power \( p_g \), reducing the ripple in the DC-link voltage \( u_{dc} \). On the other hand, the currents become unbalanced according to the instantaneous power theory and their peak value increases close to 1 p.u.

V. CONCLUSION

This paper has presented an augmented adaptive observer for grid-voltage sensorless control of a grid-connected converter equipped with an LCL filter. Linearized dynamics have been derived for tuning of the observer, and the observer has been experimentally tested. The results indicate fast tracking of the estimated quantities. The proposed observer can simultaneously estimate the positive- and negative-sequence components of the grid voltage even in highly unbalanced conditions. Grid-voltage sensorless control provides redundancy in the case of sensor faults or cost savings when the voltage sensors can be eliminated. The proposed observer could be applied, e.g., in active-front-end rectifiers of motor drives or solar inverters.

APPENDIX A

DISCRETE-TIME STATE-SPACE MODEL

The detailed expressions of the system matrix \( \Phi \) and input vector \( \Gamma_e \) in (3) have been presented in [23]. When \( u_{g+} \), \( u_{g-} \), and \( \omega_{g+} \) are assumed to be constant during the sampling period \( T_s \), the input matrices \( \Gamma_{g+} \) and \( \Gamma_{g-} \) are obtained from

\[
\Gamma_{g,m} = \left( \int_{0}^{T_s} e^{A_\Gamma (m-1)\omega (T_s - \tau)} d\tau \right) B_g
\]

where \( m = 1 \) for \( \Gamma_{g+} \) and \( m = -1 \) for \( \Gamma_{g-} \). The matrix \( A \) and the vector \( B_g \) are from the continuous-time model of the system, given in [23]. The resulting elements of the input vector \( \Gamma_{g,m} = [b_{g1,m}, b_{g2,m}, b_{g3,m}]^T \) are

\[
\begin{align*}
 b_{g1,m} &= \gamma \left[ -m \omega g+ \omega_p \sin(\omega_p T_s) + j m^2 \omega g+^2 \cos(\omega_p T_s) \\
 &- j \omega_p^2 \exp(j m \omega g+ T_s) - j \delta_m \right]/\left[ m \delta_m \omega g+ (L_{fc} + L_{fg}) \right] \\
 b_{g2,m} &= \gamma \left[ \omega_p \cos(\omega_p T_s) - \omega g+ \exp(j m \omega g+ T_s) \\
 &+ j m \omega g+ \sin(\omega_p T_s) \right]/\left[ m \delta_m \omega g+ C \right] \\
 b_{g3,m} &= \gamma \left[ m \omega g+ \omega_p L_{fc} \sin(\omega_p T_s) - j m^2 \omega g+^2 L_{fc} \cos(\omega_p T_s) \\
 &+ j \exp(j m \omega g+ T_s) \left( \delta_m L_{fg} + m^2 \omega g+^2 L_{fc} \right) - j \delta_m L_{fg} \right]/\left[ m \delta_m \omega g+ L_{fg} \right]
\end{align*}
\]

where \( \omega_p \) is the resonance frequency of the LCL filter, \( \gamma = \exp(-j \omega g+ T_s) \), and \( \delta_m = m^2 \omega g+^2 - \omega_p^2 \).

APPENDIX B

SMALL-SIGNAL LINEARIZATION OF THE ESTIMATION-ERROR DYNAMICS

The estimation-error dynamics (12) are described by the nonlinear function

\[
\dot{x}_s(k+1) = f(\dot{x}_s(k), x_s(k), u_s(k), u_{g+}(k), \ \dot{u}_{g+}(k), \dot{u}_{g+}(k), \dot{u}_{g+}(k))
\]

Equilibrium-point quantities of this system are marked with the subscript 0. If the parameters inside the observer matrices \( \Phi_s, \Gamma_{ca} \), and \( \Gamma_{ga} \) are correct (\( L_{fc} = L_{g+}, C = C \), and \( L_{fg} = L_{fg} \)) and the frequency estimation error is zero (\( \omega_{g+} = \omega_{g+} - \omega_{g0} = 0 \)), the system and observer matrices...
are equal $\Phi_{a0} = \hat{\Phi}_{a0}$, $\Gamma_{ca0} = \hat{\Gamma}_{ca0}$, $\Gamma_{ga0} = \hat{\Gamma}_{ga0}$. It follows that the nonlinear system has an equilibrium point $\{x_{a0} = 0, u_{g+0} = 0, \omega_{g+0} = 0, \bar{\vartheta}_{g+0} = 0\}$. In the vicinity of the equilibrium point, the small-signal deviation is marked with $\delta$, e.g., $\delta x_a = \bar{x}_a - x_{a0}$. In terms of the small-signal deviations, the estimation-error dynamics around the equilibrium point are

$$
\delta \ddot{x}_a(k + 1) = \left( \frac{\partial f}{\partial \delta x_a} \right)_0 \delta \dot{x}_a(k) + \left( \frac{\partial f}{\partial \delta x_c} \right)_0 \delta x_c(k) + \left( \frac{\partial f}{\partial \delta u_c} \right)_0 \delta u_c(k) + \left( \frac{\partial f}{\partial \delta u_{g+}} \right)_0 \delta u_{g+}(k) + \left( \frac{\partial f}{\partial \delta \omega_{g+}} \right)_0 \delta \omega_{g+}(k)
$$

(31)

when constant $\omega_{g+}$ is assumed, i.e., $\delta \omega_{g+} = 0$. The partial derivatives are evaluated at the equilibrium point, and they are

$$
\left( \frac{\partial f}{\partial \delta x_a} \right)_0 = \left[ \begin{array}{c} \Phi_0 \\ 0 \\ -2i \omega_{a+} + \Gamma_a C_a \end{array} \right]
$$

$$
\left( \frac{\partial f}{\partial \delta x_c} \right)_0 = \left[ \begin{array}{c} 0 \\ \frac{\partial f}{\partial \delta u_c} \right]_0 + \left( \frac{\partial f}{\partial \delta u_{g+}} \right)_0 = 0
$$

(32)

$$
\left( \frac{\partial f}{\partial \delta u_{g+}} \right)_0 = \left[ \begin{array}{c} \Gamma_{g+0} \\ 0 \\ j T_g \Gamma_{g+0} - \left( \frac{\partial f}{\partial \delta \omega_{g+}} \right)_0 \right] \times x_{a0} + \left[ \begin{array}{c} 0 \\ 0 \\ j T_g \Gamma_{g+0} - \left( \frac{\partial f}{\partial \delta \omega_{g+}} \right)_0 \end{array} \right] u_{g+0} = \Gamma_{\omega}
$$


It is to be noted that, the elements of the linearized input vector $\Gamma_{\omega}$ are generally time-varying. The time dependence originates from the rotating negative-sequence component $(u_{g+,0} = [0, 0, 0, 1] x_{a0})$ in $x_{a0}$. In order to shorten notation, the symbol $\delta$ denoting the small-signal deviations is dropped out in the linearized dynamics in Section III.

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**REFERENCES**


