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The new cosmic microwave background (CMB) temperature maps from Planck provide the highest-quality full-sky view of the surface of last scattering available to date. This allows us to detect possible departures from the standard model of a globally homogeneous and isotropic cosmology on the largest scales. We search for correlations induced by a possible non-trivial topology with a fundamental domain intersecting, or nearly intersecting, the last scattering surface (at comoving distance $\chi_{\text{rec}}$), both via a direct search for matched circular patterns at the intersections and by an optimal likelihood search for specific topologies. For the latter we consider flat spaces with cubic toroidal (T3) equal-sided chimney (T2) and slab (T1) topologies, three multi-connected spaces of constant positive curvature (dodecahedral, truncated cube and octahedral) and three spherical cases listed above, coming from the matched-circles search, is $\chi^2$ for the T3 cubic torus; $\chi^2$ for the truncated cube; and $\chi^2$ for the octahedral.

For the T3 cubic torus; $\chi^2$ for the truncated cube; and $\chi^2$ for the octahedral.

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{$\chi^2$} for the matched-circles search; $\chi^2$ for the T3 cubic torus; $\chi^2$ for the truncated cube; and $\chi^2$ for the octahedral.
1. Introduction

This paper, one of a set of papers associated with the 2013 release of data from the Planck mission (Planck Collaboration I 2014), describes the use of Planck data to limit departures from the global isotropy and homogeneity of spacetime. We will use Planck’s measurements of the cosmic microwave background (CMB) to assess the properties of anisotropic geometries (i.e., Bianchi models) and non-trivial topologies (e.g., the torus). The simplest models of spacetime are globally isotropic and simply connected. Although both are supported by both local observations and previous CMB observations, without a fundamental theory of the birth of the Universe, observational constraints on departures from global isotropy are necessary. General Relativity itself places no restrictions upon the topology of the Universe, as was recognised very early on (e.g., De Sitter 1917); most proposed theories of quantum gravity predict topology-change in the early Universe which could be visible at large scales today.

The Einstein field equations relate local properties of the curvature to the matter content in spacetime. By themselves they do not restrict the global properties of the space, allowing a universe with a given local geometry to have various global topologies. Friedmann-Robertson-Walker (FRW) models of the universe observed to have the same average local properties everywhere still have freedom to describe quite different spaces at large scales. Perhaps the most remarkable possibility is that a vanishing or negative local curvature ($\Omega_\text{0} \equiv 1 - \Omega_\text{tot} \geq 0$) does not necessarily mean that our Universe is infinite. Indeed we can still be living in a universe of finite volume due to the global topological multi-connectivity of space, even if described by the flat or hyperbolic FRW solutions. In particular, quantum fluctuations can produce compact spaces of constant curvature, both flat (e.g., Zeldovich & Starobinskii 1984) and curved (e.g., Coule & Martin 2000; Linde 2004), within the inflationary scenario.

The primary CMB anisotropy alone is incapable of constraining curvature due to the well-known geometrical degeneracy which produces identical small-scale fluctuations when the recombination sound speed, initial fluctuations, and comoving distance to the last scattering surface are kept constant (e.g., Bond et al. 1997; Zaldarriaga & Seljak 1997; Stompor & Starobinskii 1984) and curved (e.g., Coule & Martin 2000; Linde 2004), within the inflationary scenario.

and developed into the structure that we observe. Studying structure on the last scattering surface is the best-known way to probe the global organisation of our Universe and the CMB provides the most detailed and best understood dataset for this purpose.

We can also relax assumptions about the global structure of spacetime by allowing anisotropy about each point in the Universe. This yields more general solutions to Einstein’s field equations, leading to the so-called Bianchi cosmologies. For small anisotropy, as demanded by current observations, linear perturbation about the standard FRW model may be applied. A universal shear and rotation induce a characteristic subdominant deterministic signature in the CMB, which is embedded in the usual stochastic anisotropies. The deterministic CMB temperature fluctuations that result in the homogenous Bianchi models were first examined by Collins & Hawking (1973) and Barrow et al. (1985) (and subsequently Barrow 1986), however no dark energy component was included as it was not considered plausible at the time. More recently, Jaffe et al. (2006c), and independently Bridges et al. (2007), extended these solutions for the open and flat Bianchi VIIh models to include cosmologies with dark energy. It is these solutions to Bianchi VIIh models that we study in the current article. More accurate solutions were since derived by Pontzen & Challinor (2007), Pontzen (2009) and Pontzen & Challinor (2011), where recombination is treated in a more sophisticated manner and reionisation is supported. Furthermore, we note that in these works (Pontzen & Challinor 2007, 2011; Pontzen 2009) the induced CMB polarisation contributions that arise in Bianchi models have also been derived, although here focus is given to temperature contributions.

In this paper, we will explicitly consider models of global topology and anisotropy. In a chaotic inflation scenario, however, our post-inflationary patch might exhibit large-scale local topological features (“handles” and “holes”) the can mimic a global multiply-connected topology in our observable volume. Similarly, it might also have residual shear or rotation which could mimic the properties of a global Bianchi spacetime.

Planck’s ability to discriminate and remove large-scale astrophysical foregrounds (Planck Collaboration XII 2014) reduces the systematic error budget associated with measurements of the CMB sky significantly, Planck data therefore allow refined limits on the scale of the topology and the presence of anisotropy. Moreover, previous work in this field has been done by a wide variety of authors using a wide variety of data (e.g., COBE, WMAP 1-year, 3-year, 5-year, etc.) and in this work we perform a coherent analysis.

In Sect. 2, we discuss previous attempts to limit the topology and global isotropy of the Universe. In Sect. 3 we discuss the signals induced in topologically non-trivial and Bianchi universes. In Sect. 4 the Planck data we use in the analysis are presented, and in Sect. 5 the methods we have developed to detect those signals are discussed. We apply those methods in Sect. 6 and discuss the results in Sect. 7.

2. Previous results

The first searches for non-trivial topology on cosmic scales looked for repeated patterns or individual objects in the distribution of galaxies (Sokolov & Shvartsman 1974; Fang & Sato 1983; Fagundes & Wichoski 1987; Lehoucq et al. 1996; Roukema 1996; Weatherley et al. 2003; Fuji & Yoshii 2011). The last scattering surface from which the CMB is released represents the most distant source of photons in the Universe, and hence the largest scales with which we could probe the topology of the Universe. This first became possible with the
DMR instrument on the COBE satellite (Bennett et al. 1996): various searches found no evidence for non-trivial topologies (e.g., Starobinsky 1993; Sokolov 1993; Stevens et al. 1993; De Oliveira-Costa & Smoot 1995; Levin et al. 1998; Bond et al. 1998, 2000b; Rocha et al. 2004; but see also Roukema 2000b,a), but sparked the creation of robust statistical tools, along with greater care in the enumeration of the possible topologies for a given geometry (see, for example, Lachieze-Rey & Luminet 1995 and Levin 2002 for reviews). With data from the WMAP satellite (Jarosik et al. 2011), these theoretical and observational tools were applied to a high-quality dataset for the first time. Luminet et al. (2003) and Caillerie et al. (2007) claimed the low value of the low multipoles (compared to standard ΛCDM cosmology) as evidence for missing large-scale power as predicted in a closed universe with a small fundamental domain (see also Aurich 1999; Aurich et al. 2004, 2005, 2006, 2008; Aurich & Lüstig 2013; Lew & Roukema 2008; Roukema et al. 2008). However, searches in pixel space (Cornish et al. 2004; Key et al. 2007; Niauchou et al. 2004; Bielewicz & Riazuelo 2009; Dineen et al. 2005) and in harmonic space (Kunz et al. 2006) determined that this was an unlikely explanation for the low power. Bond et al. (1998, 2000a) and Riazuelo et al. (2004a,b) presented some of the mathematical formalism for the computation of the correlations induced by topology in a form suitable for use in cosmological calculations. Phillips & Kogut (2006) presented efficient algorithms for the computation of the correlation structure of the flat torus and applied it via a Bayesian formalism to the WMAP data; similar computations for a wider range of geometries were performed by Niauchou & Jaffe (2007).

These calculations used a variety of different vintages of the COBE and WMAP data, as well as a variety of different sky cuts (including the unmasked internal linear combination (ILC) map, not originally intended for cosmological studies). Nonetheless, none of the pixel-space calculations which took advantage of the full correlation structure induced by the topology found evidence for a multiply-connected topology with a fundamental domain within or intersecting the last scattering surface. Hence in this paper we will attempt to corroborate this earlier work and put the calculations on a consistent footing.

The open and flat Bianchi type VIIh models have been compared previously to both the COBE (Bunn et al. 1996; Kogut et al. 1997) and WMAP (Jaffe et al. 2005, 2006b) data, albeit with smaller sky coverage. In general, the limits on the global rotation and shear of the Universe. A statistically significant correlation between one of the Bianchi VIIh models and the WMAP ILC map (Bennett et al. 2003) was first detected by Jaffe et al. (2005). However, it was noted that the parameters of this model are inconsistent with standard constraints. Nevertheless, when the WMAP ILC map was “corrected” for the best-fit Bianchi template, some of the so-called “anomalies” reported in WMAP data disappear (Jaffe et al. 2005, 2006b; Cayón et al. 2006; McEwen et al. 2006). A modified template fitting technique was performed by Land & Magueijo (2006) and, although a statistically significant template fit was not reported, the corresponding “corrected” WMAP data were again free of many large scale “anomalies”. Subsequently, Ghosh et al. (2007) used the bipolar power spectrum of WMAP data to constrain the amplitude of any Bianchi component in the CMB. Due to the renewed interest in Bianchi models, solutions to the CMB temperature fluctuations induced in Bianchi VIIh models when incorporating dark energy were since derived by Jaffe et al. (2006c) and Bridges et al. (2007). Nevertheless, the cosmological parameters of the Bianchi template embedded in WMAP data in this setting remain inconsistent with constraints from the CMB alone (Jaffe et al. 2006a,c). Furthermore, Pontzen & Challinor (2007) compared the polarisation power spectra of the best-fit Bianchi VIIh model found by Jaffe et al. (2006a) with the WMAP 3-year data (Page et al. 2007) and also concluded that the model could be ruled out since it produced greater polarization than observed in the WMAP data. A Bayesian analysis of Bianchi VIIh models was performed by Bridges et al. (2007) using WMAP ILC data to explore the joint cosmological and Bianchi parameter space via Markov chain Monte Carlo sampling, where it was again determined that the parameters of the resulting Bianchi cosmology were inconsistent with standard constraints. In a following study by Bridges et al. (2008) it was suggested that the CMB “cold spot” (Vielva et al. 2004; Cruz et al. 2006; Vielva 2010) could be driving evidence for a Bianchi component. Recently, this Bayesian analysis has been revisited by McEwen et al. (2013) to handle partial-sky observations and to use nested sampling methods (Skilling 2004; Feroz & Hobson 2008; Feroz et al. 2009). McEwen et al. (2013) conclude that WMAP 9-year temperature data do not favour Bianchi VIIh cosmologies over ΛCDM.

3. CMB correlations in anisotropic and multiply-connected universes

3.1. Topology

All FRW models can describe multi-connected universes. In the case of flat space, there are a finite number of compactifications, the simplest of which are those of the torus. All of them have continuous parameters that describe the length of periodicity in some or all directions (e.g., Riazuelo et al. 2004b). In a space of constant non-zero curvature the situation is notably different – the presence of a length scale (the curvature radius R0) precludes topological compactification at an arbitrary scale. The size of the space must now reflect its curvature, linking topological properties to Ωm ≠ 1 − Ωk. In the case of hyperbolic spacetimes, the list of possible compact spaces of constant negative curvature is still infinite, but discrete (Thurston 1982), while in the positive curvature spherical space there is only a finite set of well-proportioned possibilities (i.e., those with roughly comparable sizes in all directions; there are also the countably infinite lens and prism topologies) for a multi-connected space (e.g., Gaussenmann et al. 2001; Riazuelo et al. 2004a).

The effect of topology is equivalent to considering the full simply-connected three-dimensional spatial slice of the space-time (known as the covering space) as being filled with repetitions of a shape which is finite in some or all directions (the fundamental domain) – by analogy with the two-dimensional case, we say that the fundamental domain tiles the covering space. For the flat and hyperbolic geometries, there are infinite copies of the fundamental domain; for the spherical geometry, with a finite volume, there is a finite number of tiles. Physical fields repeat their configuration in every tile, and thus can be viewed as defined on the covering space but subject to periodic boundary conditions. Topological compactification always break isotropy, and for some topologies also the local homogeneity of physical fields. Positively curved and flat spaces studied in this paper are homogeneous, however hyperbolic multi-connected spaces are never homogeneous.

The primary observable effect of a multi-connected universe is the existence of directions in which light could circumnavigate the space in cosmological time more than once, i.e., the radial distance Xrec to the surface of last scattering exceeds the size of the universe. In these cases, the surface of last scattering can
intersect the (notional) edge of a fundamental domain. At this intersection, we can view the same spacetime event from multiple directions – conversely, it appears in different directions when observed from a single point.

Thus, temperature perturbations in one direction, \( T(\hat{n}) \), become correlated with those in another direction \( T(\hat{m}) \) by an amount that differs from the usual isotropic correlation function \( C(\theta) \), where \( \theta \) denotes the angle between \( \hat{n} \) and \( \hat{m} \). Considering a pixelized map, this induces a correlation matrix \( C_{pp'} \) which depends on quantities other than the angular distance between pixels \( p \) and \( p' \). This break from statistical isotropy can therefore be used to constrain topological models. Hence, we need to calculate the pixel-space correlation matrix or its equivalent in harmonic space.

In this paper we consider the following topologies using the likelihood method: a) toroidal flat models with equal-length compactification size \( L \) in three directions, denoted \( T[L, L, L] \); b) toroidal flat models with different compactification lengths, parametrized by \( L_1, L_2, L_3 \), denoted \( T[L_1, L_2, L_3] \); c) three major types of single-action positively curved spherical manifolds with dodecahedral, truncated cubical and octahedral fundamental domains \( (P', O', T') \) compactification groups correspondingly, see Gaustad et al. (2001); and d) two sample negative curvature hyperbolic spaces, m004(−5,1) being one of the smallest known compact hyperbolic spaces as well as the relatively large v3543(2,3). Scales of fundamental domains of compactified curved spaces are fixed in the units of curvature and are summarised in Table 1, where we quote the volume \( V \), radius of the largest sphere that can be inscribed in the domain \( R_c \) (equal to the distance to the nearest face from the origin of the domain), the smallest sphere in which the domain can be inscribed \( R_a \) (equal to the distance to the farthest vertex), and the intermediate scale \( R_b \) that is taken to be the distance to the edges for spherical spaces and the “spine” distance for hyperbolic topologies. For the cubic torus with edge length \( L \), these lengths are \( R_c = L/2 \), \( R_a = \sqrt{3}L/2 \) and \( R_b = \sqrt{3}L/2 \). The ratio \( R_b/R_a \) is a good indicator of the shape of the fundamental domain. Note that when \( \chi_{rec} \) is less than \( R_c \), multiple images on large scales are not present, although the \( C_{pp'} \) correlation matrix is still modified versus the singly-connected limit. The effects of topology usually become strong when \( \chi_{rec} \) exceeds the intermediate \( R_b \); conversely, for flat and nearly-flat geometries, there are limits to the allowed topologies (Mota et al. 2011).

A much wider class of topologies is explicitly constrained using the matched circles method. As discussed in Sect. 5.1, because of computational limitations we restrict our analysis to pairs of circles centered around antipodal points, so called back-to-back circles. Thus, we can constrain all topologies predicting pairs of such circles. The strongest constraints are imposed on topologies predicting back-to-back circles in all directions i.e., all the single action manifolds, among them tori of any shape and the three spherical cases considered explicitly in the likelihood analysis. Weaker constraints are imposed on topologies with all back-to-back circles centred on a great circle of the celestial sphere such as half-turn, quarter-turn, third-turn and sixth-turn spaces, as well as Klein and chimney spaces. The statistic can also constrain the multi-connected spaces predicting one pair of antipodal matching circles such as Klein or chimney spaces with horizontal flip, vertical flip or half-turn and slab space translated without screw motion. Other topologies catalogued in Riazuelo et al. (2004b) are not constrained by this analysis: the Hantzsche-Wendt space; the chimney space with half-turn and flip; the generic slab space; the slab space with flip; spherical manifolds with double and linked action; and all the hyperbolic topologies including those two cases considered using the likelihood method.

### 3.1.1. Computing correlation matrices

The CMB temperature pixel-pixel correlation matrix is defined as the ensemble-average product of the temperature at two different pixels:

\[
C_{pp'} = \langle T_p T_{p'} \rangle .
\]

It can be calculated as a double radial integral of the ensemble average of the product of the source functions that describe the transport of photons through the universe from the last scattering surface to the observer:

\[
C_{pp'} = \int_0^{\chi_{rec}} dx \int_0^{\chi_{rec}} dx' \langle S(\chi \hat{q}_p) S(\chi' \hat{q}_{p'}) \rangle ,
\]

where \( \hat{q}_p \) and \( \hat{q}_{p'} \) are unit vectors that point at pixels \( p \) and \( p' \) on the sky, and \( \chi \) and \( \chi' \) are proper distances along radial rays pointing towards the last scattering surface.

Two techniques have been developed to compute the CMB correlation function for multiply-connected universes. In one approach, one constructs the orthonormal set of basis functions that satisfy the boundary conditions imposed by compactification (eigenfunctions of the Laplacian operator furnish such a basis), and assembles the spatial correlation function of the source \( \langle S(\chi \hat{q}_p) S(\chi' \hat{q}_{p'}) \rangle \) from such a basis (Cornish & Spergel 1999; Lehoucq et al. 2002). In the other approach, one applies the method of images to create the compactified version of \( \langle S(\chi \hat{q}_p) S(\chi' \hat{q}_{p'}) \rangle \) from the one computed on the universal covering space by resumming the latter over the images of the 3D spatial positions \( \chi \hat{q}_p \) (Bond et al. 1998, 2000a,b):

\[
\langle S(\chi \hat{q}_p) S(\chi' \hat{q}_{p'}) \rangle = \sum_{y \in \Gamma} \langle S(\chi \hat{q}_p) S(\gamma \hat{q}_{p'}) \rangle ,
\]

### Table 1. Parameters of analysed curved spaces.

<table>
<thead>
<tr>
<th>Size</th>
<th>Dodecahedral</th>
<th>Spherical</th>
<th>Octahedral</th>
<th>Hyperbolic</th>
</tr>
</thead>
<tbody>
<tr>
<td>( V/R_0 )</td>
<td>0.16</td>
<td>0.41</td>
<td>0.82</td>
<td>0.98</td>
</tr>
<tr>
<td>( R_c/R_0 )</td>
<td>0.31 (\pi/10)</td>
<td>0.39 (\pi/8)</td>
<td>0.45</td>
<td>0.54</td>
</tr>
<tr>
<td>( R_a/R_0 )</td>
<td>0.37</td>
<td>0.56</td>
<td>0.56</td>
<td>0.64</td>
</tr>
<tr>
<td>( R_b/R_0 )</td>
<td>0.40</td>
<td>0.58</td>
<td>0.79 (\pi/4)</td>
<td>0.75</td>
</tr>
</tbody>
</table>

(3)
where the superscripts c and u refer to the quantity in the multiply-connected space and its universal cover, respectively. The tilde refers to the need for sum regularization in the models with an infinite set of images, e.g., hyperbolic and flat toroidal ones. \( \Gamma \) is the discrete subgroup of motions which defines the multiply-connected space and \( \gamma(x) \) is the spatial point on the universal cover obtained by the action of the motion \( \gamma \in \Gamma \) on the point \( x \). Note that we can consider the location of one of the pixels as fixed and consider the action of \( \gamma \) on the other due to symmetry. This equation defines the action of \( \gamma \) on the source function itself, needed unless all the terms in the source function are scalar quantities (which is the case if one limits consideration to Sachs-Wolfe terms) when the action is trivial.

Both methods are general, but have practical considerations to take into account when one increases the pixel resolution. For computing \( C_{p'p} \) up to the resolution corresponding to harmonic mode \( \ell \approx 40 \) both methods have been tested and were found to work equally well. In this paper we employ both approaches.

The main effect of the compactification is that \( C_{p'p} \) is no longer a function of the angular separation between the pixels \( p \) and \( p' \) only, due to the lack of global isotropy. In harmonic space the two-point correlation function of the CMB is given by

\[
C_{\ell\ell'}^{\text{mm}} = \langle a_{\ell m} a_{\ell' m'} \rangle = C_{\ell} \delta_{\ell \ell'} \delta_{mm'}, \tag{4}
\]

where \( \delta_{\ell \ell'} \) is the Kronecker delta symbol and \( a_{\ell m} \) are the spherical harmonic coefficients of the temperature on the sky when decomposed into the spherical harmonics \( Y_{\ell m}(\hat{\mathbf{n}}) \) by

\[
T(\hat{\mathbf{n}}) = \sum_{\ell m} a_{\ell m} Y_{\ell m}(\hat{\mathbf{n}}). \tag{5}
\]

Note that the two-point correlation function \( C_{\ell\ell'}^{\text{mm}} \) is no longer diagonal, nor is it \( m \)-independent, as in an isotropic universe.

A flat universe provides an example when the eigenfunctions of the Laplacian are readily available in a set of plane waves. The topological compactification in the flat space discretizes the spectrum of the wavevector magnitudes \( k^2 \) and selects the subset of allowed directions. For example, for a toroidal universe the length of the fundamental cell needs to be an integer multiple of allowed directions. For example, for a toroidal universe the length of the fundamental cell needs to be an integer multiple of the wavelength of the modes. We therefore recover a discrete sum over modes \( k_n = (2\pi/L)n \) for \( n = (n_x, n_y, n_z) \) a triplet of integers, instead of an integral over \( k \),

\[
C_{\ell\ell'}^{\text{mm}} \propto \int d^3k \Delta_k(k, \Delta \eta) \Delta_k(k, \Delta \eta) P(k) \rightarrow \sum_n \Delta_n(k, \Delta \eta) \Delta_n(k, \Delta \eta) P(k_n) Y_{\ell m}(\hat{\mathbf{n}}) Y_{\ell' m'}(\hat{\mathbf{n}}), \tag{6}
\]

where \( \Delta_n(k, \Delta \eta) \) is the radiation transfer function (e.g., Bond & Efstathiou 1987; Seljak & Zaldarriaga 1996). We refer to the cubic torus with three equal sides as the T3 topology; it is also possible for the fundamental domain to be compact in only two spatial dimensions (e.g., the so-called T2 “chimney” space) or one (the T1 “slab”, similar to the “lens” spaces available in manifolds with constant positive curvature) in which case the sum is replaced by an integral in those directions. These models serve as approximations to modifications to the local topology of the global manifold (albeit on cosmological scales): for example, the chimney space can mimic a “handle” connecting different regions of an approximately flat manifold.

In Fig. 1 we show rows of the pixel-space correlation matrix for a number of multiply-connected topologies as maps, showing the magnitude of the correlation within a particular pixel. For the simply-connected case, the map simply shows the same information as the correlation function \( C(\theta) \); for the topologically non-trivial cases, we see the correlations depend on distance and direction and differ from pixel to pixel (i.e., from row to row of the matrix). In Fig. 2 we show example maps of CMB anisotropies in universes with these topologies, created by direct realisations of Gaussian fields with the correlation matrices of Fig. 1.

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![Fig. 1. Top row: correlation structure (i.e., a single row of the correlation matrix) of a simply-connected universe with isotropic correlations. For subsequent rows, the left and middle column show positively curved multiply-connected spaces (left: dedocaheiral, middle: octahedral) and the right column shows equal sided tori. The upper row of three maps corresponds to the case when the size of the fundamental domain is of the size of the diameter to the last scattering surface and hence the first evidence for large angle excess correlation appears. Subsequent rows correspond to decreasing fundamental domain size with respect to the last scattering diameter, with parameters roughly chosen to maintain the same ratio between the models.](image-url)
of a deterministic Bianchi contribution and the usual stochastic contribution that arises in the ΛCDM model. The deterministic CMB temperature fluctuations that result in the Bianchi models were derived by Barrow et al. (1985), although no dark energy component was included. More recently, Jaffe et al. (2006c), and independently Bridges et al. (2007), extended these solutions for the open and flat Bianchi VIIh models to include cosmologies with dark energy. We defer the details of the CMB temperature fluctuations induced in Bianchi models to these works and give only a brief description here.

Bianchi VIIh models describe a universe with overall rotation, parameterized by an angular velocity, ω, and a three-dimensional rate of shear, parameterized by the tensor σij; we take these to be relative to the z axis. The model has a free parameter, first identified by Collins & Hawking (1973), describing the comoving length-scale over which the principal axes of shear and rotation change orientation. The ratio of this length scale to the present Hubble radius is typically denoted x, which defines the h parameter of type VIIh models through (Barrow et al. 1985)

\[ x = \sqrt{\frac{h}{1 - \Omega_{\text{tot}}}}, \tag{7} \]

where the total energy density \( \Omega_{\text{tot}} = \Omega_m + \Omega_x \). The parameter x acts to change the “tightness” of the spiral-type CMB temperature contributions that arise due to the geodesic focusing of Bianchi VIIh cosmologies. The shear modes \( \sigma_{ij} \) of combinations of orthogonal coordinate axes are also required to describe a Bianchi cosmology. The present dimensionless vorticity \( (\omega/H)_0 \) may be related to the dimensionless shear modes \( (\sigma_{ij}/H)_0 \) by (Barrow et al. 1985)

\[ \frac{\omega}{H} = \frac{(1+h)^{1/2}(1+9h)^{1/2}}{6x^2\Omega_{\text{tot}}} \sqrt{\frac{(\sigma_{12})^2}{H}_0 + \frac{(\sigma_{13})^2}{H}_0}, \tag{8} \]

where \( H \) is the Hubble parameter. The spherical harmonic coefficients of the Bianchi VIIh induced temperature component are proportional to \( (\sigma_{12} \pm i\sigma_{13})/H_0 \) and are non-zero for azimuthal modes \( m = \pm 1 \) only (Barrow et al. 1985; McEwen et al. 2006; Pontzen & Challinor 2007). Hence, varying the phase of \( \sigma_{12} + i\sigma_{13} \) corresponds to an azimuthal rotation, i.e., a change of coordinates, while the rotationally invariant part depends on \( \sigma_{12} - i\sigma_{13} \), and we are thus free to choose equality of shear modes \( \sigma = \sigma_{12} = -\sigma_{13} \) (Pontzen & Challinor 2007), which we do for consistency with previous studies (e.g. Jaffe et al. 2005). The amplitude of the deterministic CMB temperature fluctuations induced in Bianchi VIIh cosmologies may be characterised by either \( (\sigma/H)_0 \) or \( (\omega/H)_0 \) since these parameters influence the amplitude of the induced temperature contribution only and not its morphology. The handedness of the coordinate system is also free in Bianchi VIIh models, hence both left- and right-handed models arise. Since the Bianchi-induced temperature fluctuations are anisotropic on the sky the orientation of the resulting map may vary also, introducing three additional degrees-of-freedom. The orientation of the map is described by the Euler angles\(^4\) \((\alpha, \beta, \gamma)\), where for \((\alpha, \beta, \gamma) = (0^\circ, 0^\circ, 0^\circ)\) the swirl pattern typical of Bianchi templates is centred on the South pole.

Examples of simulated Bianchi VIIh CMB temperature maps are illustrated in Fig. 3 for a range of parameters. In the analysis performed herein the BIANCHI2\(^5\) (McEwen et al. 2013)

\( \text{Fig. 3. Simulated deterministic CMB temperature contributions in Bianchi VIIh cosmologies for varying } x \text{ and } \Omega_{\text{tot}} \text{ (left-to-right } \Omega_{\text{tot}} \in \{0.10, 0.50, 0.95\}; \text{ top-to-bottom } x \in \{0.1, 0.3, 0.7, 1.5, 6.0\}) \). In these maps the swirl pattern typical of Bianchi-induced temperature fluctuations is rotated from the South pole to the Galactic centre for illustrative purposes.

\[ \text{4. Data description} \]

We use Planck maps that have been processed by the various component-separation pipelines described in Planck Collaboration XII (2014). The methods produce largely consistent maps of the sky, with detailed differences in pixel intensity, noise properties, and masks. Here, we consider maps produced by the Commander–Ruler, NILC, SMICA and SEVEM methods. Each provides its own mask and we also consider the conservative common mask.

We note that because our methods rely on rather intensive pixel- or harmonic-space calculations, in particular considering a full set of three-dimensional orientations and, for the likelihood methods, manipulation of an anisotropic correlation matrix, computational efficiency requires the use of data degraded from the native HEALPix (Górski et al. 2005) \( N_{\text{side}} = 2048 \) resolution of the Planck maps. Because the signatures of either a multiply-connected topology or a Bianchi model are most prominent on large angular scales, this does not result in a significant loss of ability to detect and discriminate amongst the models (see Sect. 5.3). However, it is worth pointing out that the likelihood and matched-circles methods are sensitive to different angular scales as applied to Planck data here. The likelihood method explicitly retains only low-\( l \) (large-scale) information in its correlation matrix, whereas the matched-circles method considers anisotropies at angular scales down to tens of arcminutes (still large in comparison to the native resolution of the Planck maps). Of course, the matched-circles method exploits the correlation of the small-scale patterns along matched circles potentially separated by large angles; this effect is not generated by

\[ \text{http://www.jasonmcewen.org/} \]

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intrinsicely large angular scale anisotropies but by the bound-
ary conditions of the fundamental domain imposed by the mul-
tiply connected topology. As described in Sect. 5.1, the matched
circles statistic used here damps the anisotropies at the largest
angular scales relative to those at smaller scales so sensitivity of
the method does not rely on the former.

The topology analyses both rely on degraded maps and
masks. The matched-circles method smooths with a 30' Gaus-
sian filter and degrades the maps to \(N_{\text{side}} = 512\), and uses a
mask derived from the SEVEM component separation method.
Because the performance of the matched-circles statistic de-
dpends on anisotropies on smaller angular scales, it can be sig-
nificantly degraded by the point source cut. As there are more
point sources detected in the Planck maps than in the WMAP
maps, the problem of point source masking is more severe in
the present case. We mask only those point sources from the
full-resolution \(f_{\text{sky}} = 0.73\) SEVEM mask with amplitude, after
smoothing and extrapolation to the 143 or 217 GHz channels,
greater than the faintest source originally detected at those fre-
nuencies. The mask derived in this way retains \(f_{\text{sky}} = 0.76\) of the
sky.

The likelihood method smooths the maps and masks with an
11° Gaussian filter and then degrades them to \(N_{\text{side}} = 16\) and conser-
vatively masks out any pixel with more than 10% of its ori-
ignal subpixels masked. At full resolution, the common mask
retains a fraction \(f_{\text{sky}} = 0.73\) of the sky, and \(f_{\text{sky}} = 0.78\) when
degraded to \(N_{\text{side}} = 16\) (the high-resolution point-source masks
are largely filled in the degraded masks). The Bianchi analysis
is performed in harmonic space, and so does not require explicit
degradation in pixel space. Rather, a noisy mask is added in pixel
space to effectively marginalise the pixel values in the masked
region (as described in more detail below; see also McEwen
et al. 2013), before the data are transformed at full resolution
into harmonic space and considered only up to a specified max-
imum harmonic \(\ell\), where correlations due to the mask are taken
into account.

Different combinations of these maps and masks are used to
discriminate between the topological and anisotropic models
described in Sect. 3.

5. Methods

5.1. Topology: circles in the sky

The first set of methods, exemplified by the circles-in-the-sky
of Cornish et al. (1998), involves a frequentist analysis using a
statistic which is expected to differ between the models exam-
nined. For the circles, this uses the fact that the intersection of
the topological fundamental domain with the surface of last scat-
tering is a circle, which one potentially views from two differ-
dent directions in a multiply-connected universe. Of course, the
matches are not exact due to noise, foregrounds, the integrated
Sachs-Wolfe (ISW) and Doppler effects along the different lines
of sight.

By creating a statistic based on the matching of differ-
ent such circles, we can compare Monte Carlo simulations of
both a simply-connected, isotropic null model with specific
anisotropic or topological models. We may then calibrate de-
tections and non-detections using Monte Carlo simulations. In
principle, these simulations should take into account the com-
lications of noise, foreground contributions, systematics, the
ISW and Doppler effects. However, they do not include gravi-
tational lensing of the CMB as the lensing deflection angle is
small compared to the minimal angular scale taken into account
in our analysis. Note that the null test is generic (i.e., not tied
to a specific topology) but any detection must be calibrated with
specific simulations for a chosen topology or anisotropic model.
A very similar technique can be used for polarisation by taking
into account the fact that the polarisation pattern itself is now
not directly repeated, but rather that the underlying quadrupole
radiation field around each point on the sky is now seen from
different directions (Bielewicz et al. 2012). These methods have
been applied successfully to COBE DMR and WMAP data, and
have recently been shown to be feasible for application to Planck
data (Bielewicz et al. 2012).

The idea of using the matched circles to study topology is
due to Cornish et al. (1998). In that work, a statistical tool was
developed to detect correlated circles in all sky maps of the CMB
anisotropy – the circle comparison statistic. In our studies we
will use version of this statistic optimised for the small-scale
anisotropies as defined by Cornish et al. (2004):

\[
S^+_{ij}(\alpha, \phi) = \frac{2}{\sum_n |m|} \left| \langle \Delta T_{ij} \rangle_m e^{-im\varphi} \right|,
\]

where \(\Delta T_{ij} \) and \(\Delta T_{ij}^+ \) denote the Fourier coefficients of the tem-
perature fluctuations around two circles of angular radius \(\alpha\) cen-
tered at different points on the sky, \(i\) and \(j\), respectively, with re-
relative phase \(\varphi\). The \(m\)th harmonic of the temperature anisotropies
around the circle is weighted by the factor \(|m|\), taking into ac-
count the number of degrees of freedom per mode. Such weight-
ing enhances the contribution of small-scale structure relative to
large-scale fluctuations and is especially important since the
large-scale fluctuations are dominated by the ISW effect. This
can obscure the image of the last scattering surface and reduce
the ability to recognise possible matched patterns on it.

The above \(S^+\) statistic corresponds to pair of circles with the
points ordered in a clockwise direction (phased). For alternative
ordering, when one of the circles the points are ordered in an
anti-clockwise direction (anti-phased), the Fourier coeffi-
cients \(\Delta T_{ij} \) are complex conjugated, defining the \(S^-\) statistic.
This allows the detection of both orientable and non-orientable
topologies. For orientable topologies the matched circles have
anti-phased correlations while for non-orientable topologies they
have a mixture of anti-phased and phased correlations.

The statistic has a range over the interval \([-1, 1]\). Circles that
are perfectly matched have \(S = 1\), while uncorrelated circles will
have a mean value of \(S = 0\). Although the statistic can also take
negative values for the temperature anisotropy generated by the
Doppler term (Bielewicz et al. 2012), anticorrelated circles are
not expected for the total temperature anisotropy considered in
this work. To find matched circles for each radius \(\alpha\), the maxi-
mum value \(\max_{\alpha} S^+_{ij}(\alpha, \phi) \) is determined.

Because general searches for matched circles are computa-
tionally very intensive, we restrict our analysis to a search for
pairs of circles centered around antipodal points, so called back-
to-back circles. As described above, the maps were also down-
graded to \(N_{\text{side}} = 512\), which greatly speeds up the computations
required, but with no significant loss of discriminatory power, as
seen in Sect. 5.3.1. More details on the numerical implementa-
tion of the algorithm can be found in the paper by Bielewicz &
Banday (2011).

As mentioned in Sect. 3.1, the constraints we will derive con-
cern topologies that predict matching pairs of back-to-back cir-
cles. However, the constraints do not apply to those universes for
which the orientation of the matched circles is impossible to de-
tect due to partial masking on the sky. Because of the larger sky
fraction removed by the Planck common mask than for WMAP
this probability is larger for the analysis of the Planck maps. Moreover, the smaller fraction of the sky used in the search of matched circles results in a false detection level larger with our \( f_{\text{sky}} = 0.76 \) mask than for the \( f_{\text{sky}} = 0.78 \) 7-year WMAP mask. As a result we obtain weaker – but more conservative – constraints on topology than for similar analyses of WMAP data (Bielawicz & Banday 2011).

To draw any conclusions from an analysis based on the statistic \( S^*_\text{max}(\alpha) \), it is very important to correctly estimate the threshold for a statistically significant match of circle pairs. We used 300 Monte Carlo simulations of CMB maps, described in detail in Sect. 5.3.1, to establish the threshold such that fewer than 1% of simulations would yield a false event.

5.2. Bayesian analyses

The second set of methods take advantage of the fact that the underlying small-scale physics is unchanged in both anisotropic and topological models compared to the standard cosmology, and thus a Gaussian likelihood function will still describe the statistics of temperature and polarization on the sky, albeit no longer with isotropic correlations. When considering specific topologies, these likelihood methods instead calculate the pixel-pixel correlation matrix. This has been done for various torus topologies (which are a continuous family of possibilities) in the flat Universe as well as for locally hyperbolic and spherical geometries (which have a discrete set of possibilities for a given value of the curvature). More general likelihood-based techniques have been developed for generic mild anisotropies in the initial power spectrum (Hanson & Lewis 2009), which may have extension to other models. For the Bianchi setting, an isotropic zero-mean Gaussian likelihood is recovered by subtracting a deterministic Bianchi component from the data, where the cosmological covariance matrix remains diagonal in harmonic space but masking introduces non-diagonal structure that must be taken into account.

Because these methods use the likelihood function directly, they can take advantage of any detailed noise correlation information that is available, including any correlations induced by the foreground-removal process. We denote the data by the vector \( \mathbf{d} \), which may be in the form of harmonic coefficients \( d_m \) or pixel temperatures \( d_p \), or, in general, coefficients of the temperature expansion in any set of basis functions. We denote the model under examination by the discrete parameter \( M \), which can take on the appropriate value denoting the usual isotropic case, or the Bianchi case, or one of the possible multiply-connected universes. The continuous parameters of model \( M \) are given by the vector \( \Theta \), which for this case we can partition into \( \Theta_C \) for the cosmological parameters shared with the usual isotropic and simply-connected case, and \( \Theta_T \) which denotes the parameters for the appropriate anisotropic case, or be it a topologically non-trivial universe or a Bianchi model. Note that all of the anisotropic cases contain “nuisance parameters” which give the orientation of either the fundamental domain or the Bianchi template which we can marginalize over as appropriate.

Given this notation, the posterior distribution for the parameters of a particular model, \( M \), is given by Bayes’ theorem:

\[
P(\Theta|\mathbf{d}, M) = \frac{P(\mathbf{d}|\Theta, M)P(\Theta|M)}{P(\mathbf{d}|M)}. \tag{10}
\]

Here, \( P(\Theta|M) = P(\Theta_C, \Theta_T|M) \) is the joint prior probability of the standard cosmological parameters \( \Theta_C \) and those describing the anisotropic universe \( \Theta_T \). \( P(\mathbf{d}|\Theta, M) \equiv \mathcal{L} \) is the likelihood, and the normalizing constant \( P(\mathbf{d}|M) \) is the Bayesian evidence, which can be used to compare the models to one another.

We will usually take the priors to be simple “non-informative” distributions (e.g., uniform over the sphere for orientations, uniform in length for topology scales, etc.) as appropriate. The form of the likelihood function will depend on the anisotropic model: for multiply-connected models, the topology induces anisotropic correlations, whereas for the Bianchi model, there is a deterministic template, which depends on the Bianchi parameters, in addition to the standard isotropic cosmological perturbations. We will assume that any other non-Gaussian signal (either from noise or cosmology) is negligible (Planck Collaboration XXIII 2014; Planck Collaboration XXIV 2014) and use an appropriate multivariate Gaussian likelihood.

Given the signal and noise correlations, and a possible Bianchi template, the procedure is similar to that used in standard cosmological-parameter estimation, with a few complications. Firstly, the evaluation of the likelihood function is computationally expensive and usually limited to large angular scales. This means that in practice the effect of the topology on the likelihood is usually only calculated on those large scales. Secondly, the orientation of the fundamental domain or Bianchi template requires searching (or marginalizing) over three additional parameters, the Euler angles.

5.2.1. Topology

In topological studies, the parameters of the model consist of \( \Theta_C \), the set of cosmological parameters for the fiducial best-fit flat cosmological model, and \( \Theta_T \), the topological parameters which include the set of compactification lengths \( L_x, L_y, L_z \) for flat toroidal model or the curvature parameter \( \Omega_K \) for curved spaces, and a choice of compactification \( T \). In our studies we keep \( \Theta_C \) fixed, and vary \( \Theta_T \) for a select choice of compactifications listed in Sect. 3.1. These parameters define the predicted two-point signal correlation matrix \( C_{pp} \) for each model, which are precomputed. Additional internal parameters, including the amplitude of the signal \( A \) and the angles of orientation of the fundamental domain of the compact space relative to the sky \( \varphi \) (e.g., parameterized by a vector of the three Euler angles), are maximized and/or marginalized over during likelihood evaluation.

The likelihood, i.e., the probability to find a temperature data map \( \mathbf{d} \) with associated noise matrix \( N \) given a certain topological model is then given by

\[
P(d|C[\Theta_C, \Theta_T, T], A, \varphi) \propto \frac{1}{\sqrt{|AC + N|}} \times \exp \left\{ -\frac{1}{2} d^\top (AC + N)^{-1} d \right\}. \tag{11}
\]

Working with a cut-sky, it is often easier to start the analysis with data and a correlation matrix given in pixel space. However, especially in the realistic case of negligible noise on large scales, the matrix \( C + N \) is poorly conditioned in pixel space, and pixel space evaluation of the likelihood is, as a rule, not robust. Indeed, there are typically more pixels than independent modes that carry information about the signal (e.g., even in the standard isotropic case, sub-arcminute pixels would not be useful due to beam-smoothing; with anisotropic correlations and masked regions of the sky, more complicated linear combinations of pixels even on large scales may have very little signal content). Therefore in general we expand the temperature map \( d_p \), the theoretical correlation matrix \( C_{pp} \) and the noise covariance matrix \( N_{pp} \) in a discrete set of mode functions \( \phi_n(p) \),

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orthonormal over the pixelized sphere, possibly with weights $w(p)$, obtaining the coefficients of expansion

$$d_n = \sum_p d_n^{p}\psi_n(p)u(p);$$

$$C_{nm} = \sum_{p} \sum_{p'} C_{pp'}\psi_n(p)\psi_{n'}(p')w(p)w(p');$$

$$N_{nm} = \sum_{p} \sum_{p'} C_{pp'}\psi_n(p)\psi_{n'}(p')w(p)w(p').$$

(12)

Next we select $N_m$ such modes for comparison and consider the likelihood marginalized over the remainder of the modes

$$p(d|\{\mathcal{C}(\Theta_C, \mathcal{T}), \varphi, A\}) \propto \frac{1}{\sqrt{|AC + N|}} \times \exp\left\{-\frac{1}{2} \sum_{n=1}^{N} d_n^{p'}(AC + N)^{-1}d_n\right\},$$

(13)

where $C$ and $N$ are restricted to the $N_m \times N_m$ block of chosen modes. Flexibility in choosing mode functions and their number $N_m$ is used to achieve the compromise between the robustness of the projected $C_{pp'}$ matrix on the one hand, and the amount of discriminating information retained in the data on the other. The weights $w(p)$ can be used to improve the accuracy of transforms on a pixelized sky.

For full-sky analysis the natural choice of the mode functions is the set of ordinary spherical harmonics $Y_{nm}(p)$ which leads to standard harmonic analysis with modes limited to a suitably chosen $\ell_{\text{max}}$. Here, where we focus on masked data, we have made a somewhat different choice. As a mode set for comparison we use the $N_{m0} = 837$ largest eigenvectors of the $C_{pp'}$ matrix, restricted to the masked sky, for the fiducial flat isotropic model with best-fit parameters $\Theta_C$. We emphasize that the correlation matrix computed for this reduced dataset has fewer modes, but contains no additional assumptions beyond those of the original $C_{pp'}$.

Since computation of $C_{pp'}$ matrices for a range of topological models is expensive, we do not aim to determine the full Bayesian evidence $P(d|T)$ which would require marginalization over all parameters $\Theta_C, \mathcal{T}$, an overall amplitude of the correlation matrix $A$ (proportional to the physical amplitude $c_T$ or the scalar amplitude $A_s$), and orientation (Euler angles) $\varphi$, and would in addition be sensitive to the prior probabilities assumed for the size of the fundamental domain. Instead we directly compare the likelihood along the changing set of $\mathcal{T}$ that has as its limit the flat fiducial model defined by $\Theta_C$. In case of toroidal topology such a limit is achieved by taking compactification lengths to infinity, while for curved models we vary $\Omega_k$ in comparison to the flat limit $\Omega_k = 0$. In the latter case, for the spherical spaces we change $\Omega_k$, and $H_0$ together with $\Omega_k$ to track the CMB geometrical degeneracy line in which the recombination sound speed, initial fluctuations, and comoving distance to the last scattering surface are kept constant (e.g., Bond et al. 1997; Zaldarriaga & Seljak 1997; Stompor & Efstathiou 1999), and for hyperbolic spaces we vary $\Omega_k$ while keeping $H_0$ and $\Omega_k - \Omega_m$ fixed to fiducial values. Note that hyperbolic multi-connected spaces, in contrast to tori and the single-action positive curvature manifolds considered in this paper, are not only anisotropic but also inhomogeneous. Therefore, the likelihood is expected to be dependent on the position of the observer. We do not study this dependence here.

For each parameter choice, we find the likelihood at the best orientation $\varphi$ of the topology with respect to the sky after marginalizing over the amplitude $A$ of the signal (hence, this can be considered a profile likelihood with respect to the orientation parameters). This likelihood is compared both with the fiducial model applied to the observed temperature map and with the likelihood of the topological model applied to the simulated realization of the isotropic map drawn from the fiducial model. Such a strategy is optimized for the detection of topological signatures. For non-detections, the marginalized likelihood can be a better probe of the overall power of the data to reject a non-trivial topology, and so for real data below, we also show the likelihood marginalized over the orientations $\varphi$. We estimate the marginalized likelihood from the random sample of 10000 orientations, drawn statistically uniformly on the $S^3$ sphere of unit quaternions representing rotations of the fundamental domain relative to the observed sky.

5.2.2. Bianchi

For the Bianchi analysis the posterior distribution of the parameters of model $M$ is given by Bayes’ Theorem, specified in Eq. (10), similar to the topological setting. The approach of McEwen et al. (2013) is followed, where the likelihood is made explicit in the context of fitting a deterministic Bianchi template embedded in a stochastic CMB background, defined by the power spectrum $C_l(\Theta_C)$ for a given cosmological model with parameters $\Theta_C$. The Bianchi $\mathcal{V}_\beta$ parameters are denoted $\Theta_\beta$. The corresponding likelihood is given by

$$p(d|\Theta_\beta, \Theta_C) \propto \frac{1}{\sqrt{|X(\Theta_C)|}} \exp[-\chi^2(\Theta_C, \Theta_\beta)/2],$$

(14)

where

$$\chi^2(\Theta_C, \Theta_\beta) = [d - b(\Theta_\beta)]^T X^{-1}(\Theta_C) [d - b(\Theta_\beta)]$$

(15)

and $d = \{d_{\text{m0}}\}$ and $b(\Theta_\beta) = \{b_{\text{m0}}(\Theta_\beta)\}$ are the spherical harmonic coefficients of the data and Bianchi template, respectively, considered up to the harmonic band-limit $\ell_{\text{max}}$. A band-limit of $\ell_{\text{max}} = 32$ is considered in the subsequent analysis for computational tractibility and since this is sufficient to capture the structure of the CMB temperature fluctuations induced in Bianchi $\mathcal{V}_\beta$ models in the vicinity of the best-fit model found in WMAP data (see, e.g., McEwen et al. 2006). The likelihood is computed in harmonic space where rotations of the Bianchi template can be performed efficiently (McEwen et al. 2006).

The covariance matrix $X(\Theta_C)$ depends on whether the full-sky or partial-sky masked setting is considered. In the full-sky setting $X(\Theta_C) = C(\Theta_C)$ as first considered by Bridges et al. (2007), where $C(\Theta_C)$ is the diagonal CMB covariance matrix with entries $C_l(\Theta_C)$ on the diagonal. In the case of a zero Bianchi component, Eq. (14) then reduces to the likelihood function used commonly to compute parameter estimates from the power spectrum estimated from CMB data (e.g., Verde et al. 2003). In the masked setting considered subsequently, the situation is a little more involved.

In order to handle a mask in the harmonic space analysis of Bianchi models we follow the approach of McEwen et al. (2013), where masking noise is added to the data to effectively marginalise over the pixel values of the data in the masked region. The masking noise $m$ is chosen to be zero-mean and large in the masked region of the data, and zero elsewhere. Consequently, the masking noise is anisotropic over the sky but...
may be chosen to be uncorrelated, and may thus be defined by its covariance
\[ \langle m(\omega_i)m'(\omega_j) \rangle = \delta_{ij} \sigma^2_m(\omega_i), \]  
(16)
where \( \delta_{ij} \) is Kronecker delta symbol, \( \omega_i \) denotes the angular coordinate of pixel \( i \), and the variance of the noise for pixel \( i \) is given by a constant value in the masked regions \( \sigma^2_m(\omega_i) = \Sigma^2_m \) and zero elsewhere. By synthetically adding masking noise that is much larger than the original data in the masked region of the sky, we effectively marginalise over the pixel values of the data in this region. The noisy mask introduces coupling in harmonic space that must be accounted for in the analysis. The covariance matrix of the resultant data is given by \( X(\Theta_C) = C(\Theta_C) + M \), where \( M \) is the non-diagonal mask covariance matrix:
\[ M_{\ell mm'} = \langle m_{\ell m} m'_{\ell m'} \rangle = \sum_{\omega_i} \sigma^2_m(\omega_i) Y^*_{\ell m}(\omega_i) Y_{\ell m'}(\omega_i) \Omega_i^2, \]  
(17)
and \( \Omega_i \) is the area of pixel \( i \) (see McEwen et al. 2013 for further details).

The \( \chi^2 \) of the likelihood for the Bianchi case hence differs from the topology case by the nonzero Bianchi template \( \Lambda \). They correspond to odds of 1 in 12, approximately, while a factor of \( 2 \) corresponds to odds of 1 in 150, approximately.

5.3. Simulations and Validation

5.3.1. Topology

Circles-in-the-Sky Before beginning the search for pairs of matched circles in the \( \text{Planck} \) data, we validate our algorithm using simulations of the CMB sky for a universe with 3-torus topology for which the dimension of the cubic fundamental domain is \( L = 2H_0^{-1} \), and with cosmological parameters corresponding to the \( \Lambda \text{CDM} \) model (see Komatsu et al. 2011, Table 1) determined from the 7-year WMAP results combined with the measurements of the distance from the baryon acoustic oscillations and the Hubble constant. We performed simulations computing directly the \( a_{\ell m} \) coefficients up to the multipole of order \( \ell_{\max} = 500 \) as described in Bielewicz & Banday (2011) and convolving them with the same smoothing beam profile as used for the data, i.e., a Gaussian beam with 30’ FWHM. In particular, we verified that our code is able to find all pairs of matched circles in such a map. The map with marked pairs of matched circles with radius \( \alpha \approx 24' \) and the statistic \( S_{\max}(\alpha) \) for the map are shown in Figs. 4 and 5, respectively. Note that the peak amplitudes in the statistic, corresponding to the temperature correlation for matched circles, decrease with radius of the circles. Cornish et al. (2004) noted that this is primarily caused by the Doppler term, which becomes increasingly anticorrelated for circles with radius smaller than 45’.

The intersection of the peaks in the matching statistic with the false detection level estimated for the CMB map corresponding to the simply-connected universe defines the minimum radius of the correlated circles which can be detected for this map. The height of the peak with the smallest radius seen in Fig. 5 indicates that the minimum radius is about \( \alpha_{\min} \approx 20' \).

For the Monte Carlo simulations of the CMB maps for the simply-connected universe we used the same cosmological parameters as for the multi-connected universe, i.e., corresponding to the \( \Lambda \text{CDM} \) model determined from the 7-year WMAP results. The maps were also convolved with the same beam profile as for the simulated map for the 3-torus universe and data, as well as masked with the same cut used for the analysis of data. The false
Bayesian Analysis

Because of the expense of the calculation of the correlation matrix, we wish to limit the number of three-dimensional wavevectors $k$ we consider, as well as the number of spherical harmonic modes $\ell$, and finally the number of different correlation matrices as a whole. We need to ensure that the full set of matrices $C_{\ell m}^{m'}$ that we calculate contains all of the available information on the correlations induced by the topology in a sufficiently fine-grained grid. For this purpose, we consider the Kullback-Leibler (KL) divergence as a diagnostic (see, e.g., Kunz et al. 2006, 2008, for applications of the KL divergence to topology). The KL divergence between two probability distributions $p_1(x)$ and $p_2(x)$ is given by

$$d_{KL} = \int p_1(x) \ln \frac{p_1(x)}{p_2(x)} \, dx.$$  

If the two distributions are Gaussian with correlation matrices $C_1$ and $C_2$, this expression simplifies to

$$d_{KL} = -\frac{1}{2} \left[ \ln \left| C_1 C_2^{-1} \right| + \text{Tr} \left( I - C_1 C_2^{-1} \right) \right],$$

and is thus a measure of the discrepancy between the correlation matrices. The KL divergence can be interpreted as the ensemble average of the log-likelihood-ratio $\Delta \ln \mathcal{L}$ between realizations of the two distributions. Hence, they enable us to probe the ability to tell if, on average, we can distinguish realizations of $p_1$ from a fixed $p_2$ without having to perform a brute-force Monte Carlo integration. Thus, the KL divergence is related to ensemble averages of the likelihood-ratio plots that we present for simulations (Fig. 10) and real data (Sect. 6), but does not depend on simulated or real data.

We first use the KL divergence to determine the size of the fundamental domain which we can consider to be equivalent to the simply-connected case (i.e., the limit in which all dimensions of the fundamental domain go to infinity). We note that in our standard $\Lambda$CDM model, the distance to the surface of last scattering is $x_{\text{rec}} \approx 3.1416(\theta_0)^{-1}$. We would naively expect that as long as the sphere enclosing the last scattering surface can be enclosed by the fundamental domain ($L = 2x_{\text{rec}}$), we would no longer see the effects of non-trivial topology. However, because the correlation matrix includes the full three-dimensional correlation information (not merely the purely geometrical effects of completely correlated points) we would see some long-scale correlation effects even for larger fundamental domains. In Fig. 6 we show the KL divergence (as a function of $(LH_0)^{-1}$ so that the simply-connected limit $L \rightarrow \infty$ is at a finite position) for the $T[L, L, L]$ (cubic), $T[L, L, 7]$ (chimney) and $T[L, 7, 7]$ (cubical) and $T[L, 7, 7]$ (slab) spaces and show that it begins to level off for $(LH_0)^{-1} \leq 1/5$, although these topologies are still distinguishable from the $T[7, 7, 7]$ torus which is yet closer to the value for a simply-connected universe $d_{KL}[7, 7, 7] = 1.1$. These figures indicate that a length of $L = 7H_0^{-1}$ is an acceptable proxy for the simply-connected infinite Universe. The figures, as well as the likelihoods computed on simulations and data, show steps and other structures on a variety of scales corresponding to the crossing of the different length scales of the fundamental domain $\mathcal{R}_0$, $\mathcal{R}_m$, and $\mathcal{R}_l$ crossing the last scattering surface; smaller fundamental domains with longer intersections with the last scattering surface are easier to detect.

Computational limitations further prevent us from calculating the likelihood at arbitrary values of the fundamental domain size parameters. We must therefore ensure that our coarse-grained correlation matrices are sufficient to detect a topology even if it lies between our gridpoints. In Fig. 7 we show the KL divergence as a function of the size of the fundamental domain, relative to various models, both aligned with our grid $(LH_0 = 4.5)$ and in between our grid points $(LH_0 = 5.25)$. We see that the peak is wide enough that we can detect a peak within $\delta LH_0 \approx 0.1$ of the correct value. We also show that we can detect anisotropic fundamental domains even when scanning through cubic tori: we show a case which approximates a “chimney” universe with one direction much larger than the distance to the last scattering surface.

Because our topological analyses do not simultaneously vary the background cosmological parameters along with those describing the topology, we also probe the sensitivity to the cosmology. In Fig. 8 we show the effect of varying the fiducial...
cosmology from the Planck Collaboration XVI (2014) best-fit values to those reported by WMAP (Komatsu et al. 2011)\(^8\). We see that this induces a small bias of \(\delta L_0 \approx 0.2\) but does not hinder the ability to detect a non-trivial topology. This indicates that small deviations from the correct background cosmology do not hinder our ability to detect (or rule out) topological signals.

We have also directly validated the topological Bayesian techniques with simulations. In Fig. 9 we show the log-likelihood for the above \(T[2,2,2]\) simulations as a function of two of the Euler angles, maximized over the third. We find a strong peak at the correct orientation, with a multiplicity due to the degenerate orientations corresponding to the faces of the cube (there are peaks at the North and South poles, which are difficult to see in this projection). Note that the peaks correspond to ratios of more than \(\exp(700)\) compared to the relatively smooth minima elsewhere.

In Fig. 10 we also test the ability of the Bayesian likelihood technique to detect the compactification of the space in the simulated temperature realizations drawn from the dodecahedral closed model. For curved geometries, the size of the fundamental domain is fixed with respect to the varying curvature scale \((R_0)\), whereas the distance to the last scattering \(\chi_{\text{rec}}\) is constant. Hence we plot the likelihood as a function of \(\chi_{\text{rec}}/R_0\), inversely proportional to the scale of the fundamental domain.

Two multiply-connected realizations of the sky were tested: one corresponding to the space in which the last scattering sphere can be just inscribed into the fundamental domain, \(\chi_{\text{rec}} = R_c\), when just the first large angle correlations appear, and the second drawn from a somewhat smaller space for which \(\chi_{\text{rec}} = R_e\). We see detections in both cases, stronger as the fundamental domain shrinks relative to \(\chi_{\text{rec}}\). We also calculate the likelihood for a model known to be simply-connected.

Note that the likelihood taken at the best orientations of the compact models generically shows a slight increase relative to that for the limiting simply connected space as one brings the size of the fundamental domain down to the size of the last scattering surface \((\chi_{\text{rec}} \approx R_i)\), followed, in the absence of signal in the map, by a rapid drop as soon as the models smaller than \(\chi_{\text{rec}}\) are applied. This small increase is also present in the fiducial exactly isotropic sky, a single realization of which is shown in the figure, but is a generic feature irrespective of the topology being tested (occurring also in models with \(R < R_i\)), and thus should not be taken as an indicator for compact topology. The reason for the increase is the possibility of aligning the model with a weak anisotropic correlation feature with chance patterns of a single sky realization. However the fit drastically worsens as soon as the correlation features in a model become pronounced. Moreover, the feature becomes considerably less significant when the likelihood is marginalized over the orientation (Euler angles) of the fundamental domain.

All of these results (KL divergences and likelihoods) were computed with \(\ell_{\text{max}} = 40\), corresponding approximately to \(N_{\text{side}} = 16\), indicating that this is more than adequate for detecting even relatively small fundamental domains such as the...
We note that it is difficult to compress the content of these likelihood figures down to limits upon the size of the fundamental domain. This arises because it is difficult to provide a physically-motivated prior distribution for quantities related to the size of the fundamental domain. Most naive priors would diverge towards arbitrarily large fundamental domain sizes or would otherwise depend on arbitrary limits to the topological parameters.

5.3.2. Bianchi

The ANICOSMO code (McEwen et al. 2013) is used to perform a Bayesian analysis of Bianchi VII$_h$ models, which has been extensively validated by McEwen et al. (2013) already; we briefly summarise the validation performed for the masked analysis. In McEwen et al. (2013) a CMB map is simulated, in which a simulated Bianchi temperature map with a large vorticity (i.e., amplitude) is embedded, before applying a beam, adding isotropic noise and applying a mask. Both the underlying cosmological and Bianchi parameters used to generate the simulations are well recovered. For this simulation the coupled Bianchi model is favoured over ΛCDM, with a log-Bayes factor of $\Delta \ln E \sim 50$. As expected, one finds that the log-Bayes factor favours ΛCDM in simulations where no Bianchi component is added. For further details see McEwen et al. (2013).

6. Results

We now discuss the results of applying the circles-in-the-sky and likelihood methods to Planck data to study topology and Bianchi VII$_h$ cosmologies.

6.1. Topology

Neither the circles-in-the-sky search nor the likelihood method find evidence for a multiply-connected topology. We show the matched circle statistic in Fig. 11. We do not find any statistically significant correlation of circle pairs in any map. As seen in Fig. 5, the minimum radius at which the peaks expected for the matching statistic are larger than the false detection level is $r_{\text{min}} \approx 20''$. Thus, we can exclude at the confidence level of 99% any topology that predicts matching pairs of back-to-back circles larger than this radius, assuming that relative orientation of the fundamental domain and mask allows its detection. This implies that in a flat universe described otherwise by the Planck fiducial ΛCDM model, a 99% confidence-limit lower bound on the size of the fundamental domain is $L/2 \gtrsim \chi_{\text{rec}} \cos(\alpha_{\text{min}}) = 0.94 \chi_{\text{rec}} = 13.2$ Gpc. This is better than the limits from the marginalized likelihood ratios below for the tori and octahedron topologies and slightly worse than the limits for the dodecahedron and truncated cube. However, this constraint is not limited only to these few topologies. The frequentist analysis provides constraints upon a much wider class of topologies than those explicitly considered in the Bayesian likelihood approach; it concerns all topologies listed in Sect. 3.1.

The likelihood method also show no evidence of a multiply-connected universe. We present the likelihood for various models. In Fig. 12 we show the likelihood (marginalized over amplitude and maximized over orientation of the fundamental domain) for the cubic torus, fixing the background cosmology to the best-fit flat Universe Planck model (Planck Collaboration XVI 2014). We see that this is maximized for $L > 2\chi_{\text{rec}}$, i.e., showing no evidence for non-trivial topology.
Note that the likelihood shows mild features as the size goes through the other scales associated with the topology, in particular a small increase in the likelihood when the scale of the inscribed sphere \( R \) is crossed. However, the same increase is found when the toroidal model is compared to a single realization of a strictly isotropic fiducial sky, and thus, should not be interpreted as a detection of multi-connected topology. The origin of this likelihood behaviour at best fit angles is that the freedom of orientation can be used to align small enhancements in large-angle correlations in the anisotropic \( L = 2R \) model with random features in the given single realization of the sky. When marginalized over all possible orientations the effect is significantly reduced; the slight rise is \( \Delta \ln L \approx 1.9 \) from a likelihood of \( P = 650 \), which is comparable to the numerical noise inherent in our stochastic integration. For even smaller spaces, more extensive correlations of the temperature can no longer be accommodated and for \( L < 2R_m \) the likelihood of the T3 cubic toroidal model drops quickly, although not strictly monotonically.

In Figs. 13 and 14 we show the likelihood for the \( T[L, L, 7] \) chimney and \( T[L, 7, 7] \) slab topologies, which are also maximized in the simply-connected limit. The T2 chimney, with only two compact dimensions, is less constrained than the T3 cube, and the T1 slab, with one compact dimension, even less so.

We find similar limits for the topologies allowed in a closed universe with a locally spherical geometry. In Fig. 15 we show the likelihood for the dodecahedral fundamental domain, in which the likelihood for the dodecahedral fundamental domain, in
Fig. 13. Same as Fig. 12, but for a toroidal space with one large dimension fixed at $7H_0^{-1}$ and two short dimensions of equal size $L$ (approximating the “chimney” space). $R_i$ and $R_m$ are marked while $R_\infty = \infty$.

Fig. 14. Same as Fig. 12, but for a toroidal space with two large dimensions fixed at $7H_0^{-1}$ and one short dimension of variable $L$ (approximating the “slab” space). $R_i$ is marked while $R_m = R_\infty = \infty$.

Fig. 15. Top: likelihood as a function of the distance to last scattering surface in curvature units for a locally spherical multiply-connected universe with a dodecahedral ($I'$) fundamental domain with $R_i = 0.31R_0$. Lines are for different estimates of the CMB temperature from Planck data as in Fig. 12. In this figure, the $\chi^2_{rec}/R_i$ parameterizes the position of the model on the geometrical degeneracy line which links $H_0$ and $\Omega_k$ with $\Omega_m$. The degeneracy relations are approximated as $\Omega_k = 0.691 + 2.705\Omega_m$ and $H_0 = 67.8 + 388\Omega_m + 1200\Omega_m^2$. As in the toroidal case, there is no detection of a small space at the level expected from the simulations of Sect. 5. Fundamental domains larger than the last scattering diameter are preferred for the dodecahedral and truncated cube spaces with somewhat weaker restriction for the octahedral case. Note that an observationally motivated prior on $H_0$ or $\Omega_\Lambda$ would be yet more restrictive on the fundamental domain size. For all three topologies, again as in the toroidal case, the maximum of the likelihood at best fit orientation is detected for the finite volume spaces with $\chi^2_{rec} \approx R_i$ at the level $\Delta \ln L \approx +4$ relative to the fiducial flat simply-connected model. Since this feature is seen in the isotropic fiducial sky as well, we cannot take it as an indication of a detection of a multi-connected space. In the case of curved spaces we see that this mild increase disappears when we consider the likelihood marginalized over orientations.

We present numerical limits for these flat and positively curved spaces in Table 2. Because of the one-sided nature of these limits, we characterize the shape of the likelihood by the steepness of its fall from the value as the scale of the fundamental domain goes to infinity (i.e., the simply-connected limit). Hence, we show limits for $\Delta \ln L < -5$, (roughly equivalent to a $3\sigma - 99\%$ confidence limit – fall for a Gaussian; because of the very steep gradient, the $2\sigma$ limits are very similar) and $\Delta \ln L < -12.5$ ($5\sigma$). Note that the limits differ depending on whether we marginalize or maximize the likelihood over the orientation angles. We show lower limits on the quantity $R_i (L/2$ for a torus with edge length $L$) in units of the last scattering distance $\chi^\text{rec}$ (in conventional units, $\chi^\text{rec} \approx 14 \text{Gpc}$ for the fiducial Planck parameters; Planck Collaboration XVI 2014). In most cases, the limits are roughly $R_i \geq \chi^\text{rec} - \text{the scale of the fundamental domain must be greater than that of the last scattering surface. We place the most restrictive limits on the dodecahedron with } R_i > 1.03\chi^\text{rec} \text{ using marginalized values for } \Delta \ln L < -5.$
Conversely, the chimney and slab spaces are less constrained as the expected correlations are weaker in one or two directions; for the slab space, we only constrain $R_i = L/2 > 0.5\chi_{rec}$.

In Fig. 18 we show the likelihood for the two hyperbolic models listed in Table 1, which also show no detection of the multi-connected topology. In the hyperbolic case we space the range of space sizes by varying $\Omega_L$ while keeping $\Omega_m = \Omega_\Lambda$ as well as $H_0$ constant at fiducial values.

All of these results show at least some increase in the likelihood for certain orientations when one of the characteristic scales of the fundamental domain ($R_m$, $R_i$, or $R$) just exceed the surface of last scattering, and so no longer produces matched patterns, but induces extra correlations at large angular separations. Chance patterns can then mimic these correlations, and this is exacerbated by our conservative sky masks, which allow arbitrary patterns in the masked regions.

### 6.2. Bianchi

Masked Planck data are analysed for evidence of a Bianchi VII$_h$ component, where the prior parameter ranges adopted are the same as those specified by McEwen et al. (2013). The analysis is performed on the SMICA component-separated map, using the mask defined for this method, and is repeated on the SEVEM component-separated map for validation purposes (using the mask defined for the SEVEM method). The Bayes factors for the various Bianchi VII$_h$ models and the equivalent standard cosmological models are shown in Table 3.

For the phenomenological flat-decoupled-Bianchi model, evidence in support of a left-handed Bianchi template is found. On the Jeffreys scale (Jeffreys 1961), evidence for this model would be referred to as strong for the SMICA map and significant for the SEVEM map. For both SMICA and SEVEM component-separated data, recovered posterior distributions for the flat-decoupled-Bianchi model are shown in Fig. 19a, where similar
posterior distributions are recovered for both component separation methods. Recall that the Bianchi parameters are decoupled from the standard cosmology in the flat-decoupled-Bianchi model, hence for this model $Ω_m^B$ and $Ω_Λ^B$ are specific to the Bianchi model and should not be compared with standard values. The maximum a posteriori (MAP) best-fit template found for SMICA component-separated data is shown in Fig. 20b, with the difference between this template and the template found in WMAP 9-year data (McEwen et al. 2013) shown in Fig. 21. Note that the template found in Planck data is very similar to the template found in WMAP 9-year data (McEwen et al. 2013), which in turn is similar to the template first found by Jaffe et al. (2005). However, the template found in WMAP 9-year data (McEwen et al. 2013) is only significant in full-sky data, but not when the 9-year KQ75 WMAP mask (Bennett et al. 2013) is applied. Since the Planck SMICA and SEVEM masks are less conservative than the KQ75 mask, these findings suggest data near the Galactic plane may be playing a considerable role in supporting a Bianchi component in Planck data. The SMICA CMB map and a Bianchi-subtracted version of this map are also shown in Fig. 20. The best-fit parameters of the templates found in Planck SMICA and SEVEM component-separated data are displayed in Table 4, for both the MAP and mean-posterior estimates. The analysis was also performed on a SMICA component-separated Gaussian simulation, yielding a null detection (i.e., no evidence for a Bianchi component), as expected.

For the most physically motivated open-coupled-Bianchi model where the Bianchi VII$\text{h}$ model is coupled to the standard cosmology, there is no evidence in support of a Bianchi contribution. Recovered posterior distributions for the open-coupled-Bianchi model are shown in Fig. 19b for both SMICA and SEVEM component-separated data. Although the cosmological Bianchi parameters agree reasonably well between these different component-separated data, the posterior distributions recovered for the Euler angles differ. For SEVEM data, an additional mode of the posterior distribution is found; the mode found with SMICA data is still present in SEVEM data but is not dominant. Consequently, the best-fit estimates for the Euler angles differ between the SMICA and SEVEM component-separated data. Note that the additional mode found in SEVEM data is also present in WMAP 9-year data (McEwen et al. 2013). The resulting best-fit parameters for the open-coupled-Bianchi model are displayed in Table 5, while the corresponding MAP best-fit maps are shown in Fig. 22. Nevertheless, for both SMICA and SEVEM data the Bayes factors computed (Table 3) do not favour the inclusion of any Bianchi component for the open-coupled-Bianchi model. Planck data thus do not provide evidence in support of Bianchi VII$\text{h}$ cosmologies. However, neither is it possible to conclusively discount Bianchi VII$\text{h}$ cosmologies in favour of $\Lambda$CDM cosmologies. The constraints $(ω/H_0)^2 < 7.6 \times 10^{-10}$ (95% confidence level) on the vorticity of the physical coupled Bianchi VII$\text{h}$ left-handed models and $(ω/H_0)^2 < 8.1 \times 10^{-10}$ (95% confidence level) for right-handed models are recovered from SMICA component-separated data.

### Table 3. Log-Bayes factor relative to equivalent $\Lambda$CDM model (positive favours Bianchi model).

<table>
<thead>
<tr>
<th>Model</th>
<th>SMICA</th>
<th>SEVEM</th>
</tr>
</thead>
<tbody>
<tr>
<td>Flat-decoupled-Bianchi (left-handed)</td>
<td>$2.8 \pm 0.1$</td>
<td>$1.5 \pm 0.1$</td>
</tr>
<tr>
<td>Flat-decoupled-Bianchi (right-handed)</td>
<td>$0.5 \pm 0.1$</td>
<td>$0.5 \pm 0.1$</td>
</tr>
<tr>
<td>Open-coupled-Bianchi (left-handed)</td>
<td>$0.0 \pm 0.1$</td>
<td>$0.0 \pm 0.1$</td>
</tr>
<tr>
<td>Open-coupled-Bianchi (right-handed)</td>
<td>$-0.4 \pm 0.1$</td>
<td>$-0.4 \pm 0.1$</td>
</tr>
</tbody>
</table>

7. Discussion

We have used the Planck temperature anisotropy maps to probe the large-scale structure of spacetime. We have calculated the Bayesian likelihood for specific topological models in universes with locally flat, hyperbolic and spherical geometries, all of which find no evidence for a multiply-connected topology with a fundamental domain within the last scattering surface. After calibration on simulations, direct searches for matching circles resulting from the intersection of the fundamental topological domain with the surface of last scattering also give a null result at high confidence. These results use conservative masks of the sky, unlike previous WMAP results, which used full-sky internal linear combination maps (not originally intended for cosmological studies) or less conservative foreground masks. Hence, the results presented here, while corroborating the previous non-detections, use a single, self-consistent, and conservative dataset. The masked sky also increases the possibility of chance patterns in the actual sky mimicking the correlations expected for topologies with a characteristic scale near that of the last scattering surface.

Depending on the shape of the fundamental domain, we find $R_\frac{\chi_{\text{rec}}}{\chi}$ (Table 2) with detailed 99% confidence limits (considering the likelihood marginalized over the orientation of the fundamental domain) varying from $0.9\chi_{\text{rec}}$ for the cubic.
torus in a flat universe to $1.03\chi_{\text{rec}}$ for the dodecahedron in a positively curved universe, with somewhat weaker constraints for poorly-proportioned spaces that are considerably larger along some directions. In the case of the torus and octahedron topologies, a tighter constraint of $0.94\chi_{\text{rec}}$ comes from the matched circles method (albeit with a somewhat different interpretation of frequentist and Bayesian limits). The constraint derived using this method applies to a wide class of topologies, listed in Sect. 3.1, predicting matching pairs of back-to-back circles.

Note that the results derived using the likelihood method make use of the expected pixel-space correlations as a unique signal of non-trivial topology. Hence, although a small fundamental domain will suppress power on the largest scales of the CMB, observation of such low power on large scales as observed by COBE (Bond et al. 2000c), and confirmed by WMAP
Planck Collaboration: *Planck* 2013 results. XXVI.

**Fig. 20.** Best-fit template of left-handed flat-decoupled-Bianchi VII\(_h\) model subtracted from *Planck* SMICA component-separated data. Before subtraction, the peak-to-peak variation is ±594 \(\mu\text{K}\), reduced to ±564 \(\mu\text{K}\) after subtraction.

(Luminet et al. 2003), is not sufficient for the detection of topology. Conversely, because our methods search directly for these correlations (and indeed marginalize over the amplitude of fluctuations), a slight modification of the background FRW cosmology by lowering power in some or all multipoles (Planck Collaboration XV 2014) will not affect the ability to detect the correlations induced by such topologies.
Similarly, using a Bayesian analysis we find no evidence for a physical, anisotropic Bianchi VII$_h$ universe. However, Planck data do provide evidence supporting a phenomenological Bianchi VII$_h$ component, where the parameters of the Bianchi component are decoupled from standard cosmology. The resulting best-fit Bianchi VII$_h$ template found in Planck data is similar to that found in WMAP data previously (Jaffe et al. 2005; McEwen et al. 2013). However, although this Bianchi component can produce some of the (possibly anisotropic) temperature patterns seen on the largest angular scales (see also Planck Collaboration XXIII 2014), there is no set of cosmological parameters which can simultaneously produce these patterns and the observed anisotropies on other scales. Moreover, the parameters of the best-fit Bianchi VII$_h$ template in the decoupled setting are in strong disagreement with other measurements of the cosmological parameters.

These results are expected from previous measurements from COBE and WMAP, but Planck’s higher sensitivity and lower level of foreground contamination provides further confirmation. We have shown that the results are insensitive to the details of the preparation of the temperature maps (in particular, the method by which the cosmological signal is separated

<p>| Table 4. Parameters recovered for left-handed flat-decoupled-Bianchi model. |
|-----------------------------------|-----------------|-----------------|-----------------|</p>
<table>
<thead>
<tr>
<th><strong>Bianchi parameter</strong></th>
<th><strong>MAP</strong></th>
<th><strong>SMICA Mean</strong></th>
<th><strong>SEVEM Mean</strong></th>
</tr>
</thead>
<tbody>
<tr>
<td>$\Omega^b_0$</td>
<td>0.38</td>
<td>0.32 ± 0.12</td>
<td>0.35</td>
</tr>
<tr>
<td>$\Omega^b_{\lambda}$</td>
<td>0.20</td>
<td>0.31 ± 0.20</td>
<td>0.22</td>
</tr>
<tr>
<td>$x$</td>
<td>0.63</td>
<td>0.67 ± 0.16</td>
<td>0.66</td>
</tr>
<tr>
<td>$(\omega/H)_b$</td>
<td>$8.8 \times 10^{-10}$</td>
<td>$(7.1 \pm 1.9) \times 10^{-10}$</td>
<td>$9.4 \times 10^{-10}$</td>
</tr>
<tr>
<td>$\alpha$</td>
<td>38.1</td>
<td>51.3 ± 47.9</td>
<td>40.5</td>
</tr>
<tr>
<td>$\beta$</td>
<td>28.2</td>
<td>33.7 ± 19.7</td>
<td>28.4</td>
</tr>
<tr>
<td>$\gamma$</td>
<td>309.2</td>
<td>292.2 ± 51.9</td>
<td>317.0</td>
</tr>
</tbody>
</table>

**Notes.** Planck data favour the inclusion of a Bianchi component in this phenomenological model.

<p>| Table 5. Parameters recovered for left-handed open-coupled-Bianchi model. |
|-----------------------------------|-----------------|-----------------|-----------------|</p>
<table>
<thead>
<tr>
<th><strong>Bianchi parameter</strong></th>
<th><strong>MAP</strong></th>
<th><strong>SMICA Mean</strong></th>
<th><strong>SEVEM Mean</strong></th>
</tr>
</thead>
<tbody>
<tr>
<td>$\Omega_b$</td>
<td>0.05</td>
<td>0.07 ± 0.05</td>
<td>0.09</td>
</tr>
<tr>
<td>$\Omega^b_0$</td>
<td>0.41</td>
<td>0.33 ± 0.07</td>
<td>0.41</td>
</tr>
<tr>
<td>$\Omega^b_{\lambda}$</td>
<td>0.55</td>
<td>0.60 ± 0.07</td>
<td>0.50</td>
</tr>
<tr>
<td>$x$</td>
<td>0.46</td>
<td>0.44 ± 0.24</td>
<td>0.38</td>
</tr>
<tr>
<td>$(\omega/H)_b$</td>
<td>$5.9 \times 10^{-10}$</td>
<td>$(4.0 \pm 2.4) \times 10^{-10}$</td>
<td>$9.3 \times 10^{-10}$</td>
</tr>
<tr>
<td>$\alpha$</td>
<td>57.4</td>
<td>122.5 ± 960.0</td>
<td>264.1</td>
</tr>
<tr>
<td>$\beta$</td>
<td>54.1</td>
<td>70.8 ± 35.5</td>
<td>79.6</td>
</tr>
<tr>
<td>$\gamma$</td>
<td>202.6</td>
<td>193.5 ± 77.4</td>
<td>90.6</td>
</tr>
</tbody>
</table>

**Notes.** Planck data do not favour the inclusion of a Bianchi component in this model and some parameters are not well constrained.
from astrophysical foreground contamination). Future Planck measurement of anisotropic CMB polarization will allow us to further test models of non-trivial geometries and will provide more definitive conclusions, for example allowing us to moderately extend the sensitivity to large-scale topology (Bielewicz et al. 2012).

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